# Numerical simulation of random curves lecture 1 

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## Introduction

Theme: simulations that are related to Schramm-Loewner Evolution (SLE) in some way.

Pedagogic- Ask questions! tgk@math.arizona.edu
Goal: Enable the participants to do state of the art simulations related to SLE.

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- Simulation of SLE itself; driving function $\Rightarrow$ random curves
- Inverse problem: random curves $\Rightarrow$ driving function
- Lattice models related to SLE: LERW, percolation, SAW, Ising


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On-line at http://www.math.arizona.edu/~tgk

- These slides
- More detailed notes (pdf)
- Computer code (C++/linux)


## 2 Simulations involving conformal maps

Conformal map: from domain $D$ to domain $D^{\prime}$
Bijection that preserves angles
Bijection that is analytic function
Riemann mapping theorem If $D, D^{\prime}$ are simply connected domains, there is a conformal map between them.

Simply connected means no holes in domain
Not unique, 3 real degrees of freedom
Boundary of domain need not be smooth

### 2.1 Loewner equation - crash course

$$
\mathbb{H}=\{z: \operatorname{Im}(z)>0\}
$$

Let $\gamma$ be a curve in $\mathbb{H}$, starting at 0 , with no self intersections (simple).
$\mathbb{H} \backslash \gamma[0, t]$ is a simply connected domain.
$\Rightarrow$ conformal map $g_{t}$ from it onto $\mathbb{H}$
Simple curves $\Leftrightarrow$ one parameter families of conformal maps.
Parameterize so that

$$
g_{t}(z)=z+\frac{2 t}{z}+O\left(\frac{1}{|z|^{2}}\right), \quad z \rightarrow \infty
$$

Then Loewner's equation says

$$
\frac{\partial g_{t}(z)}{\partial t}=\frac{2}{g_{t}(z)-U_{t}}, \quad g_{0}(z)=z
$$

for some real valued function $U_{t}$ on $[0, \infty)$, the driving function.

## Loewner equation - continued

Tip of curve $\leftrightarrow$ driving function

$$
g_{t}(\gamma(t))=U_{t}, \quad \gamma(t)=g_{t}^{-1}\left(U_{t}\right)
$$

Need a limit here.
$\gamma(0)=0 \Rightarrow U_{0}=0$.
More generally, $\gamma(0)=U_{0}$.
Smoothness of $U_{t}$ : Solution to Loewner need not exist for all times $t$.
Let $K_{t}$ be the set of points $z$ in $\mathbb{H}$ for which the solution to this equation no longer exists at time $t$.

Nice $U_{t} \Rightarrow K_{t}$ is a curve $\gamma[0, t]$.
But even for a continuous $U_{t}$ it need not be.

## Stochastic Loewner equation

Random curves $\Rightarrow$ random $U_{t}$, i.e., stochastic process.
Schramm's discovery was that conformal invariance and a certain "Markov property" imply that $U_{t}$ must be a Brownian motion with mean zero.

SLE is what you get by taking $U_{t}$ to be a Brownian motion with mean zero.

For $U_{t}$ equal to Brownian motion, whether $K_{t}$ is a curve depends on the variance.

In our simulations our approximation to $U_{t}$ will be nice enough that it produces a curve.

## Loewner equation as sequence of compositions

$g_{t+s}: \mathbb{H} \backslash \gamma[0, t+s] \rightarrow \mathbb{H}$.
Do this in two stages:
$g_{s}: \mathbb{H} \backslash \gamma[0, s] \rightarrow \mathbb{H}, \quad \mathbb{H} \backslash \gamma[0, s+t] \rightarrow \mathbb{H} \backslash g_{s}(\gamma[s, s+t])$,
$\bar{g}_{t}: \mathbb{H} \backslash g_{s}(\gamma[s, t+s]) \rightarrow \mathbb{H} \quad$ (usual normalization)
Uniqueness $\Rightarrow g_{s+t}=\bar{g}_{t} \circ g_{s}, \quad$ i.e., $\quad \bar{g}_{t}=g_{s+t} \circ g_{s}^{-1}$, so

$$
\frac{d}{d t} \bar{g}_{t}(z)=\frac{d}{d t} g_{s+t}\left(g_{s}^{-1}(z)\right)=\frac{2}{g_{s+t}\left(g_{s}^{-1}(z)\right)-U_{s+t}}=\frac{2}{\bar{g}_{t}(z)-U_{s+t}}
$$

Note that $\bar{g}_{0}(z)=z$. Thus $\bar{g}_{t}(z)$ is obtained by solving the Loewner equation with driving function $\bar{U}_{t}=U_{s+t}$.
This driving function starts at $U_{s}$, and so the curve $\bar{\gamma}(t)$ associated with $\bar{g}_{t}$ starts growing at $U_{s}$.

## Loewner eq. as sequence of compositions - cont.

Partition the time interval $[0, \infty): 0=t_{0}<t_{1}<t_{2}<\cdots t_{n}<\cdots$

$$
\begin{gathered}
\bar{g}_{k}=g_{t_{k}} \circ g_{t_{k-1}}^{-1} \\
g_{t_{k}}=\bar{g}_{k} \circ \bar{g}_{k-1} \circ \bar{g}_{k-2} \circ \cdots \cdots \bar{g}_{2} \circ \bar{g}_{1}
\end{gathered}
$$

$\bar{g}_{k}(z)$ is obtained by solving the Loewner equation with driving function $U_{t_{k-1}+t}$ for $t=0$ to $t=\Delta_{k}$, where $\Delta_{k}=t_{k}-t_{k-1}$.
$\bar{g}_{k}$ maps $\mathbb{H}$ minus a curve starting at $U_{t_{k-1}}$ to $\mathbb{H}$.
Intervals of driving function $\leftrightarrow$ composition of maps
It is convenient to shift by $U_{t_{k-1}}$ so our curves start at 0 .

## Loewner eq. as sequence of compositions - cont.

Define

$$
g_{k}(z)=\bar{g}_{k}\left(z+U_{t_{k-1}}\right)-U_{t_{k-1}},
$$

It is obtained by solving the Loewner equation with driving function $U_{t_{k-1}+t}-U_{t_{k-1}}$ for $t=0$ to $t=\Delta_{k}$.
Driving function goes from 0 to $\delta_{k}=U_{t_{k}}-U_{t_{k-1}}$.
$g_{k}$ takes $\mathbb{H}$ minus a curve starting at the origin onto $\mathbb{H}$. and maps tip of that curve to $\delta_{k}$.

## The key idea

Two types of simulations:

- Given driving function, find the curve it generates.
- Given a curve, find the corresponding driving function.

For both problems the key idea is the same.
Approximate the driving function on the interval $\left[t_{k-1}, t_{k}\right]$ by a function for which the Loewner may be explicitly solved.

Maps $\bar{g}_{k}$ and $g_{k}$ can then be found explicitly.
$g_{t}$ is approximated by the composition of the appropriate maps.
Need some explicit solutions to Loewner equation

## Tilted slits

$$
f(z)=\left(z+x_{l}\right)^{1-\alpha}\left(z-x_{r}\right)^{\alpha},
$$

maps $\mathbb{H}$ to $\mathbb{H} \backslash \Gamma$ where $\Gamma$ is a line segment from 0 to a point $r e^{i \alpha \pi}$. It maps $\left[-x_{l}, x_{r}\right]$ onto $\Gamma$.

Unfortunately, its inverse cannot be explicitly computed.
For normalization we need $(1-\alpha) x_{l}=\alpha x_{r}$.

$$
f_{t}(z)=\left(z+2 \sqrt{t} \sqrt{\frac{\alpha}{1-\alpha}}\right)^{1-\alpha}\left(z-2 \sqrt{t} \sqrt{\frac{1-\alpha}{\alpha}}\right)^{\alpha}
$$

produces slit with capacity $2 t$.

## Tilted slits - continued

$g_{t}=f_{t}^{-1}$ solves Loewner eq. with driving function

$$
U_{t}=c_{\alpha} \sqrt{t}
$$

where

$$
c_{\alpha}=\frac{2(1-2 \alpha)}{\sqrt{\alpha(1-\alpha)}}
$$

NB: $\alpha \pi$ is usual polar angle. Lawler's $\alpha$ is my $1-\alpha$.

## Vertical slits

Let

$$
g_{t}(z)=\sqrt{(z-\delta)^{2}+4 t}+\delta
$$

Then it is easy to check $g_{t}$ satisfies Lowener's equation with a constant diving function: $U_{t}=\delta$.

Since the driving function does not start at 0 , the curve will not start at the origin.
The curve is just a vertical slit from $\delta$ to $\delta+2 \sqrt{t}$ i.
Vertical slits mean that we approximate the driving function by a discontinuous piecewise constant function.
Result is not a curve.

### 2.2 Simulating SLE

Recall that $g_{t}(\gamma(t))=U_{t}$. Define $z_{k}=g_{t_{k}}^{-1}\left(U_{t_{k}}\right)$. So

$$
z_{k}=\bar{g}_{k}^{-1} \circ \bar{g}_{k-1}^{-1} \circ \cdots \bar{g}_{1}^{-1}\left(U_{t_{k}}\right)
$$

Recall solving Loewner with driving function $U_{t_{k-1}+t}-U_{t_{k-1}}$ for $t=0$ to $t=\Delta_{k}$, gives $g_{k}(z)$ where $g_{k}(z)=\bar{g}_{k}\left(z+U_{t_{k-1}}\right)-U_{t_{k-1}}$. Define

$$
h_{k}(z)=g_{k}(z)-\delta_{k}=\bar{g}_{k}\left(z+U_{t_{k-1}}\right)-U_{t_{k}}
$$

where $\delta_{k}=U_{t_{k}}-U_{t_{k-1}}$. Then

$$
h_{k} \circ h_{k-1} \circ \cdots \circ h_{1}\left(z_{k}\right)=\bar{g}_{k} \circ \bar{g}_{k-1} \circ \cdots \circ \bar{g}_{1}\left(z_{k}\right)-U_{t_{k}}=0
$$

Define $f_{k}=h_{k}^{-1}$ so

$$
z_{k}=f_{1} \circ f_{2} \circ \cdots \circ f_{k}(0)
$$

## Simulating SLE - continued

$g_{k}$ maps $\mathbb{H}$ minus a curve that starts at 0 onto $\mathbb{H}$, sending the tip of the curve to $\delta_{k}$.
So $h_{k}$ maps $\mathbb{H}$ minus the curve onto $\mathbb{H}$, sending the tip to 0 .
So $f_{k}=h_{k}^{-1}$ maps $\mathbb{H}$ onto $\mathbb{H}$ minus the curve, sending 0 to the tip.
The first map $f_{k}$ welds together a small interval on $\mathbb{R}$ containing the origin to produce a small cut.
Origin is mapped to the tip of this cut.
The second map $f_{k-1}$ welds together a (possibly different) small interval to produces another small cut.

The original cut is moved away from the origin with its base being at the tip of the new cut.

This process continues. Each map introduces a new small cut whose tip is attached to the image of the base of the previous cut.

## Choice of $\delta_{k}, \Delta_{k}$

$$
\Delta_{k}=t_{k}-t_{k-1}, \quad \delta_{k}=U_{t_{k}}-U_{t_{k-1}}
$$

We approximate $U_{t}$ on each time interval $\left[t_{k-1}, t_{k}\right]$ so that $g_{k}(z)$ is known explicitly.
The two constraints on $g_{k}$ are that the curve must have capacity $2 \Delta_{k}$ and $g_{k}$ must map the tip of the curve to $\delta_{k}$.
Different choices of how we choose $g_{k}$ subject to these constraints give us different discretizations of the curve.
However, this choice will not have a significant effect.
Of much greater importance is how we choose the $\Delta_{k}$ and $\delta_{k}$.

## Uniform $\Delta_{k}$

Uniform $\Delta_{k}$ is simplest. For $\kappa$ not too large this works reasonably well Figure shows a curve with $\kappa=8 / 3$ and $N=10,000$ points.


## Uniform $\Delta_{k}$ - continued

For larger values of $\kappa$, uniform $\Delta_{k}$ is a disaster. Figure shows $\kappa=6$.


## Adaptive $\Delta_{k}$

For larger values of $\kappa$, we should use a varying $\Delta_{k}$. (Rhode).


## Capacity is not "length"

The (half plane) capacity $C$ of a set $A$ is

$$
g(z)=z+\frac{C}{z}+O\left(\frac{1}{z^{2}}\right)
$$

where $g$ maps $\mathbb{H} \backslash A$ onto $\mathbb{H}$, usual normalizations.
More intuitive def:

$$
C=\lim _{y \rightarrow \infty} y E^{i y}\left[\operatorname{Im}\left(B_{\tau}\right)\right]
$$

where $B_{t}$ is 2 d Brownian motion started at $i y$ $\tau$ is the first time it hits $A$ or $\mathbb{R}$.

## Adaptive $\Delta_{k}$

Fix a spatial scale $\epsilon$.
Start with uniform $\Delta_{k}$. Compute $z_{k}$.
Look for $z_{k}$ such that $\left|z_{k}-z_{k-1}\right| \geq \epsilon$
For these time intervals $\left[t_{k-1}, t_{k}\right]$, divide the interval into two equal intervals.

Sample Brownian motion at midpoint of $\left[t_{k-1}, t_{k}\right]$ using Brownian bridge.

Compute new $z_{k}, \cdots$.
Repeat until all $\left|z_{k}-z_{k-1}\right| \leq \epsilon$.

## Effect of choice of cut

$\kappa=8 / 3$, tilted slits vs. vertical slits, about 35,000 points


## Effect of choice of cut - continued

 $\kappa=8 / 3$, tilted slits vs. vertical slits - enlarged.

## Effect of choice of cut

$\kappa=6$, tilted slits vs. vertical slits, about 35,000 points


## Effect of choice of cut - continued

$\kappa=6$, tilted slits vs. vertical slits - enlarged.


### 2.3 Inverse SLE: curve $\Rightarrow$ driving function

Given curve, compute the driving function.
Why would you want to do this?
Usually the parameterization of the curve is not by capacity.
So if $g_{s}: \mathbb{H} \backslash \gamma[0, s] \rightarrow \mathbb{H}$ then

$$
g_{s}(z)=z+\frac{C(s)}{z}+O\left(\frac{1}{z^{2}}\right)
$$

Coefficient $C(s)$ is capacity of $\gamma[0, s]$.
The value of the driving function at $t$ is $U_{t}=g_{s}(\gamma(s)), 2 t=C(s)$.
Algorithm for finding the conformal map is "zipper" algorith.

## Shifting conformal maps - philosophy

We find it more convenient to work with the conformal map

$$
h_{s}(z)=g_{s}(z)-U_{t}
$$

where $C(s)=t$.
$h_{s}: \mathbb{H} \backslash \gamma[0, s] \rightarrow \mathbb{H}$, sends tip of $\gamma(s)$ to the origin.
The value of the driving function at $s$ is minus the constant term in the Laurent expansion of $h_{s}$ about $\infty$ :

$$
h_{s}(z)=z-U_{t}+\frac{2 C(s)}{z}+O\left(\frac{1}{z^{2}}\right),
$$

## Sequence of compositions

Let $\gamma$ be a curve in $\mathbb{H}$ starting at 0 .
Let $z_{0}, z_{1}, \cdots, z_{n}$ be points along the curve with $z_{0}=0$.
In many applications these are lattice sites.
Let $2 t_{k}$ be capacity of curve up to $z_{k}$.
Define $\bar{g}_{k}, g_{k}, h_{k}, f_{k}$ as before, so

$$
z_{k}=f_{1} \circ f_{2} \circ \cdots \cdots \circ \circ f_{k-1} \circ f_{k}(0)
$$

To simulate SLE we knew the $f_{k}$ and computed the $z_{k}$.

$$
h_{k} \circ h_{k-1} \circ \cdots h_{2} \circ h_{1}\left(z_{k}\right)=0
$$

Now we know the $z_{k}$ and want to find the $h_{k}$.

## Finding the $h_{k}$

Suppose $h_{1}, h_{2}, \cdots, h_{k}$ have been defined.
So $h_{k} \circ h_{k-1} \circ \cdots h_{2} \circ h_{1}$ sends $z_{k}$ to 0
It should send $z_{k+1}$ to a point close to origin :

$$
w_{k+1}=h_{k} \circ h_{k-1} \circ \cdots \circ h_{1}\left(z_{k+1}\right)
$$

Let $\gamma_{k+1}$ be a short simple curve that ends at $w_{k+1}$, e.g., tilted slit or vertical slit.

Let $h_{k+1}$ be corresponding conformal map.
$h_{k+1}$ has two real degrees of freedom, determined by $w_{k+1}$.

## Finding the driving function

Let $2 \Delta_{i}$ be capacity of $h_{i}$
Let $\delta_{i}$ be final value of the driving function for $h_{i}$. So

$$
h_{i}(z)=z-\delta_{i}+\frac{2 \Delta_{i}}{z}+O\left(\frac{1}{z^{2}}\right)
$$

Then

$$
h_{k} \circ h_{k-1} \circ \cdots \circ h_{1}(z)=z-U_{t_{k}}+\frac{2 t_{k}}{z}+O\left(\frac{1}{z^{2}}\right)
$$

where

$$
t_{k}=\sum_{i=1}^{k} \Delta_{i}, \quad U_{t_{k}}=\sum_{i=1}^{k} \delta_{i}
$$

Driving function of the curve is sum of driving functions of the elementary conformal maps $h_{i}$.

## Zipper algorithm for conformal maps

Digression: original zipper algorithm for computing conformal maps. See Marshall/Rhode paper.
$D$ is simply connected domain. $z_{0}, z_{1}, \cdots z_{n}$ are points on its boundary (counter clock-wise).

$$
\phi(z)=i \sqrt{\frac{z-z_{1}}{z-z_{0}}}
$$

$z_{0} \rightarrow \infty, z_{1} \rightarrow 0, \quad z_{2}, z_{3}, \cdots, z_{n}$ are mapped to curve $\gamma$ in $\mathbb{H}$.
Plane is mapped to $\mathbb{H}, \quad D$ to one side of $\gamma$ in $\mathbb{H}$.
Unzip $\gamma$. At end $D$ is mapped to quarter plane.
Map quarter plane to $\mathbb{H}$ or unit disc or ...

## Lecture 2

- Faster algorithm for simulating SLE
- Faster algorithm for computing driving function.
- Results - pictures of SLE
- Results - driving process of self-avoiding walk
- Open problems

Intermission Show movie at intermission

