

## 1. Research Goals

There are three types of errors that accrue in solving PDEs numerically:

- Model error:** Formulation of the (weak) PDE problem to be solved.
- Discretization error:** Reduction of the smooth formulation to a discrete problem.
- Solver error:** Numerical solution of the linear system.

Solving PDEs on large meshes requires mesh decimation as a pre-processing step. Decimation is guided by a **real-valued cost function on mesh edges** as follows: edge values are computed, the lowest cost edge is collapsed, and the process is iterated as needed. **Current methods** select cost functions independent of the PDE to be solved:

select cost function → decimate input → define PDE problem → select solver → solve

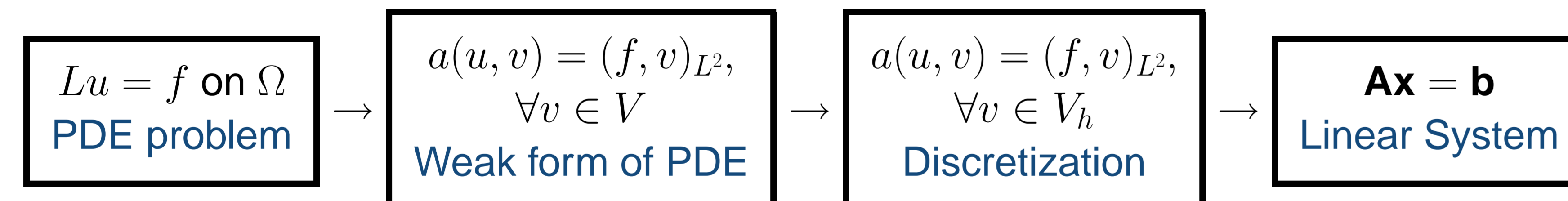
A **stable mesh decimation pipeline** must select a cost function that minimizes both discretization error and numerical error for the specific PDE problem in question:

define PDE problem → select solver → select cost function → decimate input → solve

## 2. Methodology and Background

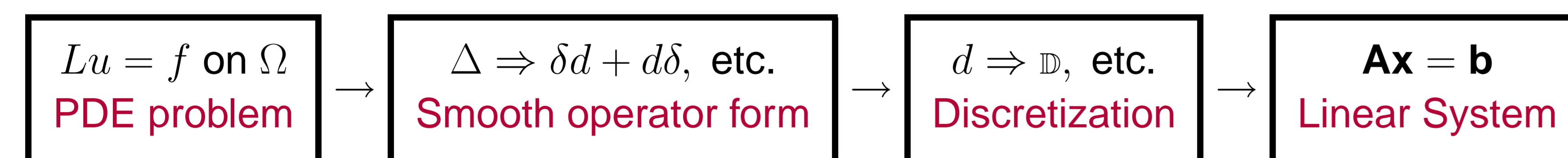
Large disparities in edge lengths (i.e. larger relative  $h$  values) cause worse stability estimates and thus greater discretization error. Large entries in the matrix of a linear system cause poor condition numbers and thus greater solver error. A stable decimation cost function should provide for **both small  $h$  values and small entries in the system matrix.**

The **Galerkin Finite Element Method** discretizes the solution space of the PDE:



Here, the entries of matrix **A** are a functional  $a(\cdot, \cdot)$  on the basis functions of  $V_h$ . Hence, **decimation must also be guided by the geometry-sensitive components of the linear functional  $a(\cdot, \cdot)$ .**

**Discrete Exterior Calculus methods** discretize the operators of the PDE:



In this case, the entries of matrix **A** depend heavily on the geometry-sensitive operators such as  $*$  and  $\delta$  and less on the topology sensitive operators such as  $d$ . Hence, **decimation must also be guided by the definition of the discrete operators.**

## Poisson-Boltzmann Electrostatics

The linear Poisson-Boltzmann (PB) equations for electrostatics are:

$$\nabla^2 \phi^{\text{int}}(r_p) = -\frac{1}{\epsilon_{\text{int}}} \sum_{k=1}^N q_k \delta(r_p - r_k), \quad p \in \Omega, \quad (1)$$

$$\nabla^2 \phi^{\text{ext}}(r_p) = \kappa^2 \phi^{\text{ext}}(r_p), \quad p \in \Omega, \quad (2)$$

A **mixed Galerkin finite element method** by Lu, Zhang, and McCammon [3] solves for surface potential  $f = \phi^{\text{int}} = \phi^{\text{ext}}$  and the normal derivative  $h = \partial \phi^{\text{ext}} / \partial n$ :

$$\begin{pmatrix} \mathbb{B} & \mathbb{A} \\ \mathbb{D} & \mathbb{C} \end{pmatrix} \begin{pmatrix} f \\ h \end{pmatrix} = \begin{pmatrix} Q \\ 0 \end{pmatrix}$$

The four sub matrices **A**, **B**, **C** and **D** have entries of the form

$$\sum_t \int_{\text{mesh facet } t} \mathcal{I} dA \quad \text{where} \quad \mathcal{I} \in \left\{ G(x_i, x_j), \frac{\partial G(x_i, x_j)}{\partial n}, u(x_i, x_j), \frac{\partial u(x_i, x_j)}{\partial n} \right\}$$

The functions  $G$  and  $u$  are the Green functions for (1) and (2), respectively:

$$G(x_i, x_j) = \frac{1}{4\pi r_{ij}} \quad \text{and} \quad u(x_i, x_j) = \frac{\exp(-\kappa r_{ij})}{4\pi r_{ij}} \quad \text{where} \quad r_{ij} = |x_i - x_j|$$

Hence, the **terms of the stiffness matrix decay like  $1/r_{ij}$**  at worst. We **use the cost function  $f_{pb}$**  to avoid small  $r_{ij}$  values, i.e. to keep Gaussian quadrature points evenly spaced.

## Generalized Born Electrostatics

The Generalized Born (GB) model of electrostatic solvation defines electrostatic solvation energy in terms of the pairwise interaction between atoms  $i$  and  $j$ :

$$G_{\text{pol}} = -\frac{\tau}{2} \sum_{i,j} \frac{q_i q_j}{[r_{ij}^2 + R_i R_j \exp(-r_{ij}^2 / 4R_i R_j)]^{1/2}} \quad \text{where} \quad \tau = \frac{1}{\epsilon_p} - \frac{1}{\epsilon_w}$$

$$\begin{aligned} \epsilon_p, \quad \epsilon_w &= \text{solute and solvent dielectric constants, respectively} \\ q_i, \quad R_i &= \text{charge and effective Born radius of atom } i, \text{ respectively} \\ r_{ij} &= \text{distance between atom } i \text{ and atom } j \end{aligned}$$

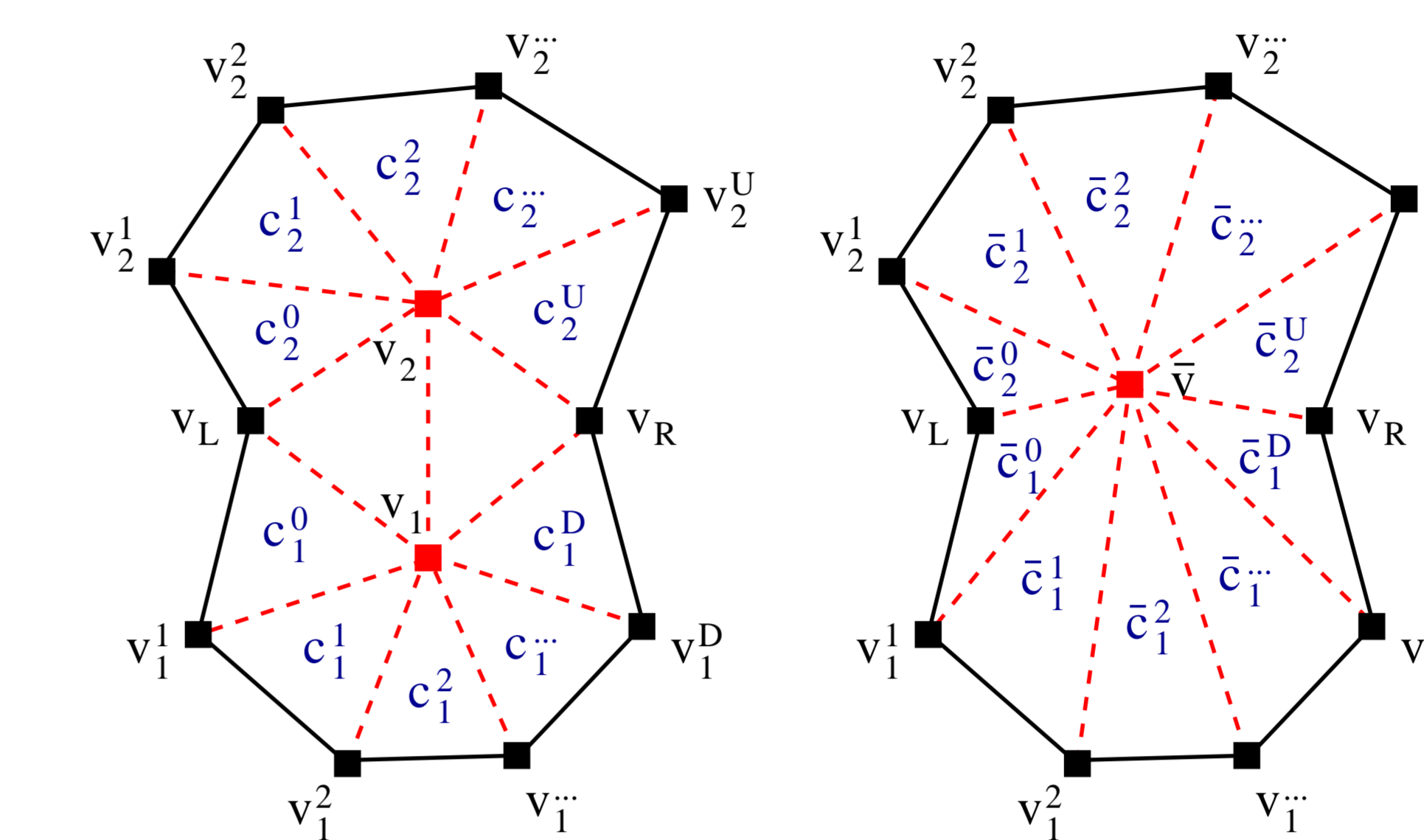
The success of a GB method hinges upon efficient and accurate approximation of the effective Born radii  $R_i$ . Bajaj and Zhao [2] use a surface integration technique to approximate

$$R_i^{-1} \approx \frac{1}{4\pi} \sum_{k=1}^N w_k \frac{(\mathbf{r}_k - \mathbf{x}_i) \cdot \mathbf{n}(\mathbf{r}_k)}{|\mathbf{r}_k - \mathbf{x}_i|^4} dS, \quad i = 1, \dots, M$$

$\mathbf{r}_k$  = Gaussian quadrature nodes with weights  $w_k$  on a triangular mesh of surface  $\Gamma$

Hence, the computation of  $R_i^{-1}$  is **sensitive to changes in  $|\mathbf{r} - \mathbf{x}_i|$** , i.e. changes in the position of Gauss points relative to the nearest atomic centers. We **use the cost function  $f_{gb}$**  to minimize the cumulative change in  $|\mathbf{r} - \mathbf{x}_i|$  values.

## 3. Problem-Specific Cost Functions



Notation for collapse of edge  $(\mathbf{v}_1, \mathbf{v}_2)$  to  $\bar{\mathbf{v}}$ . Black vertices and edges are unchanged in the collapse while red vertices and dashed edges may change. The  $c_i$  denote circumcenters:

$$c_2^u := \frac{1}{3} (\mathbf{v}_2 + \mathbf{v}_2^u + \mathbf{v}_2^{u+1}), \quad u = 0, 1, \dots, U$$

$$c_1^d := \frac{1}{3} (\mathbf{v}_1 + \mathbf{v}_1^d + \mathbf{v}_1^{d+1}), \quad d = 0, 1, \dots, D$$

**Cost Function for PB:**  $f_{pb}(\mathbf{v}_1, \mathbf{v}_2; \bar{\mathbf{v}}) := \sum_i q(T_{c_i}) - q(T_{c_i})$

•  $f_{pb}$  is larger for edge collapses which bring Gaussian quadrature points closer together.

•  $q(T)$  = quality(triangle  $T$ ) :=  $\frac{\text{longest side of } T}{\text{shortest side of } T} + \frac{\text{maximum angle of } T}{\text{minimum angle of } T}$

smaller discr. error [1] larger numerical error larger discr. error smaller numerical error

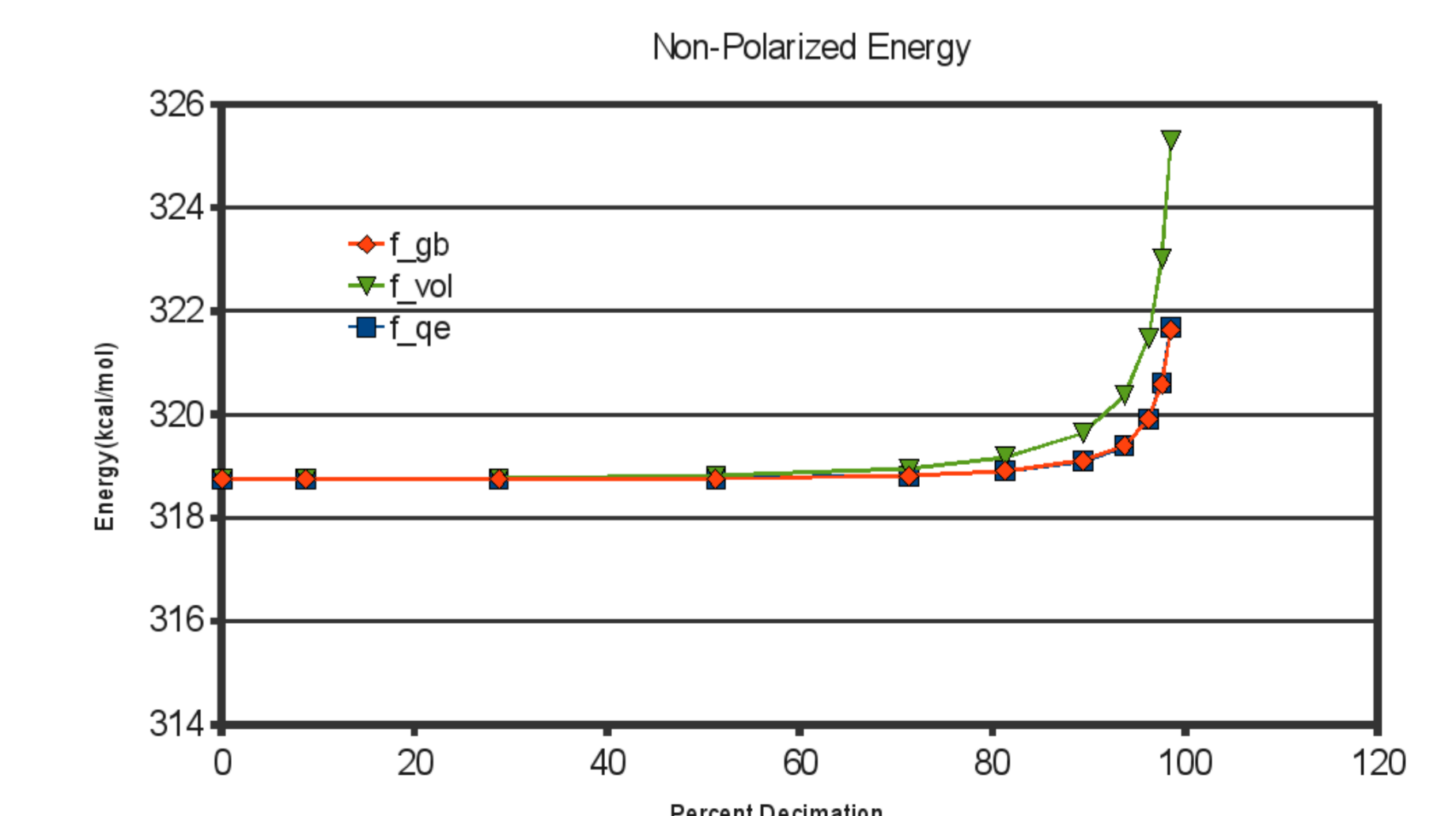
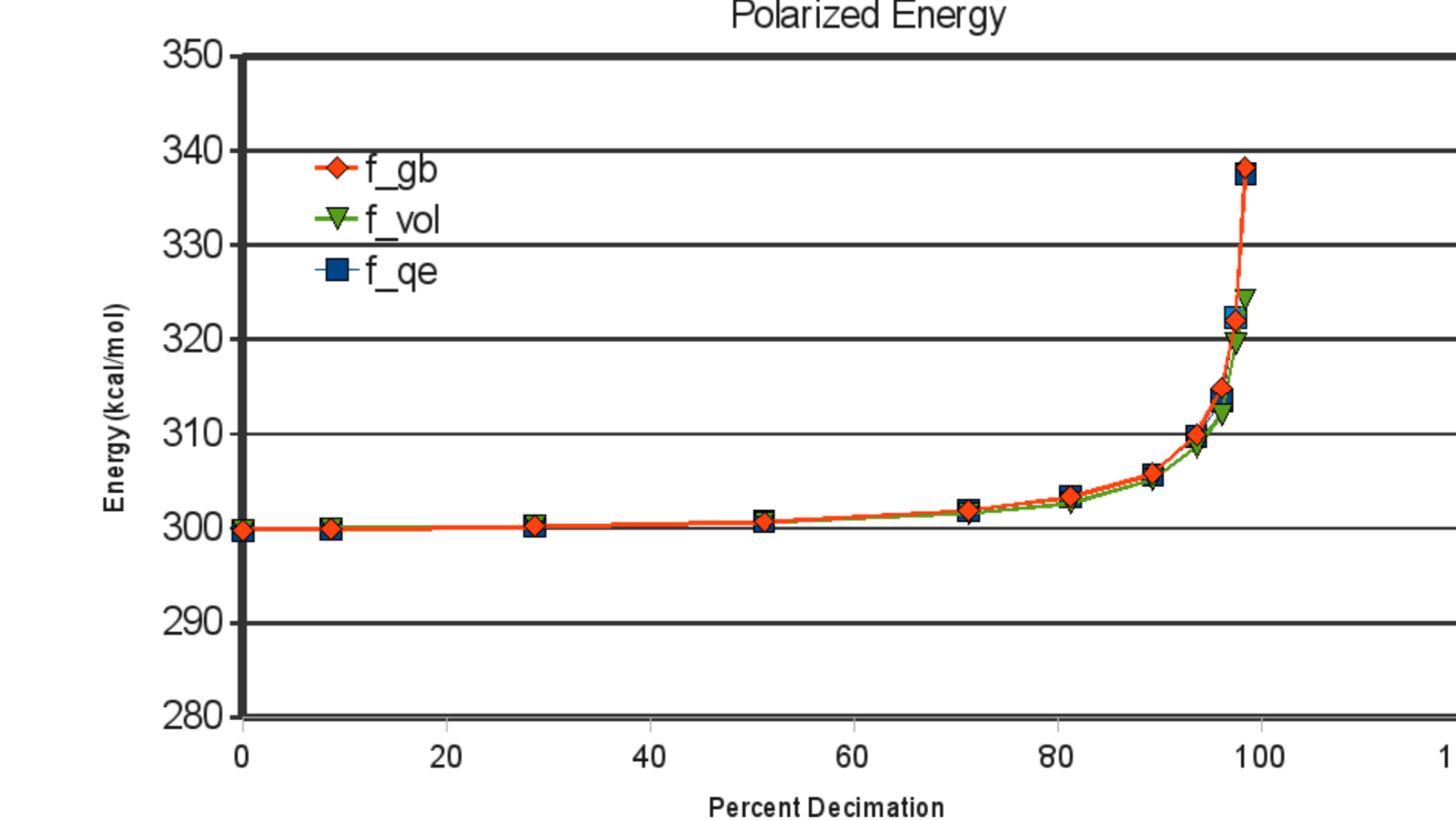
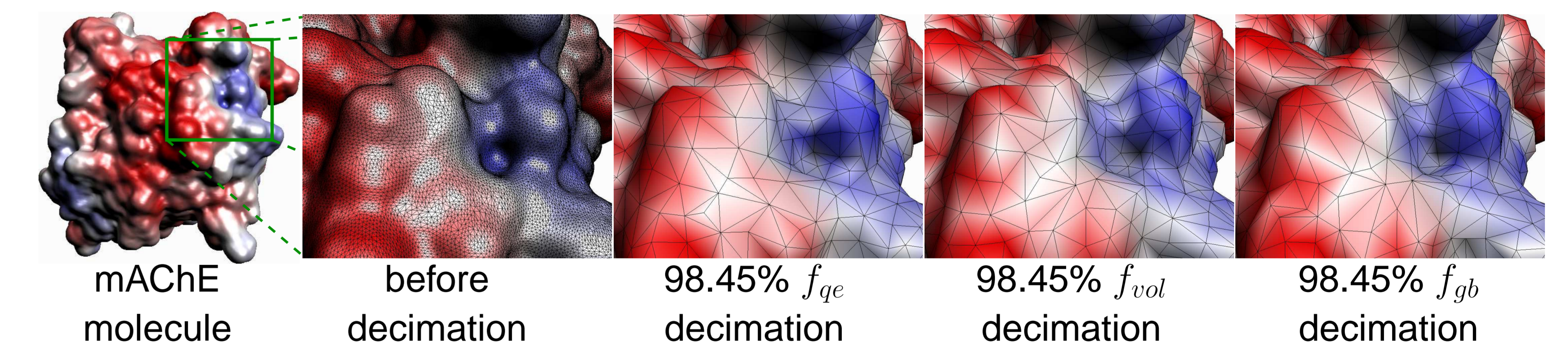
**Cost Function for GB:**  $f_{gb}(\mathbf{v}_1, \mathbf{v}_2; \bar{\mathbf{v}}) := f_{geom}(\mathbf{v}_1, \mathbf{v}_2; \bar{\mathbf{v}}) + \lambda |\sum_{i,j} |c_i - x_j|^2 - |\bar{c}_i - x_j|^2|$

•  $f_{gb}$  is larger for edge collapses with a larger cumulative change in  $|c_i - x_j|$  values.

•  $f_{geom}$  is any geometric error measure, such as  $|\mathbf{v}_1 - \mathbf{v}_2|$  or a quadratic error measure.

•  $\{x_j\}$  is the set of atomic centers lying within a fixed distance  $\rho$  (usually 2 - 5 Å) of  $\mathbf{v}_1$  or  $\mathbf{v}_2$ .

## 4. Experimental Results



## References

- [1] I. Babuska and A. K. Aziz. On the angle condition in the finite element method. *SIAM Journal on Numerical Analysis*, 13(2):214–226, 1976.
- [2] C. Bajaj and W. Zhao. Fast molecular solvation energetics and forces computation. *SIAM J. Sci. Comp.*, Submitted.
- [3] B. Lu, D. Zhang, and J. A. McCammon. Computation of electrostatic forces between solvated molecules determined by the poisson-boltzmann equation using a boundary element method. *Journal of Chemical Physics*, 122(21):214102–1–7, 2005.

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