

# Dynamics of the Elastic Pendulum

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# Agenda

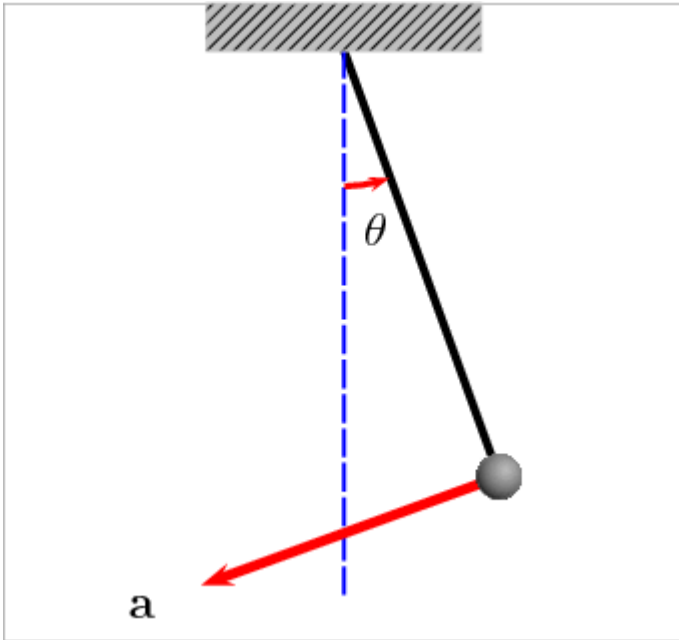
- Introduction to the elastic pendulum problem
- Derivations of the equations of motion
- Real-life examples of an elastic pendulum
- Trivial cases & equilibrium states
- MATLAB models

# The Elastic Problem (Simple Harmonic Motion)



- $F_{net} = m \frac{d^2x}{dt^2} = -kx \rightarrow \frac{d^2x}{dt^2} = -\left(\frac{k}{m}\right)x$
- Solve this differential equation to find
$$x(t) = c_1 \cos(\omega t) + c_2 \sin(\omega t) = A \cos(\omega t - \varphi)$$
- With velocity and acceleration
$$v(t) = -A\omega \sin(\omega t + \varphi)$$
$$a(t) = -A\omega^2 \cos(\omega t + \varphi)$$
- Total energy of the system
$$E = K(t) + U(t)$$
$$= \frac{1}{2}mv^2 + \frac{1}{2}kx^2 = \frac{1}{2}kA^2$$

# The Pendulum Problem (with some assumptions)



- With position vector of point mass  $\vec{x} = l (\sin\theta\vec{i} - \cos\theta\vec{j})$ , define  $\vec{r}$  such that  $\vec{x} = l\vec{r}$  and  $\vec{\theta} = \cos\theta\vec{i} + \sin\theta\vec{j}$
- Find the first and second derivatives of the position vector:

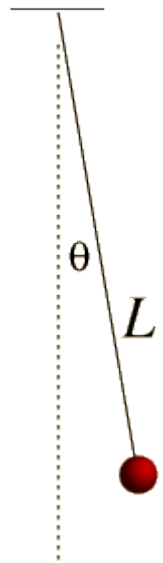
$$\frac{d\vec{x}}{dt} = l \frac{d\theta}{dt} \vec{\theta}$$

$$\frac{d^2\vec{x}}{dt^2} = l \frac{d^2\theta}{dt^2} \vec{\theta} - l \left( \frac{d\theta}{dt} \right)^2 \vec{r}$$

- From Newton's Law, (neglecting frictional force)

$$m \frac{d^2\vec{x}}{dt^2} = \vec{F}_g + \vec{F}_t$$

# The Pendulum Problem (with some assumptions)



Defining force of gravity as  $\vec{F}_g = -mg\vec{j} = mg\cos\theta\vec{r} - mg\sin\theta\vec{\theta}$  and tension of the string as  $\vec{F}_t = -T\vec{r}$ :

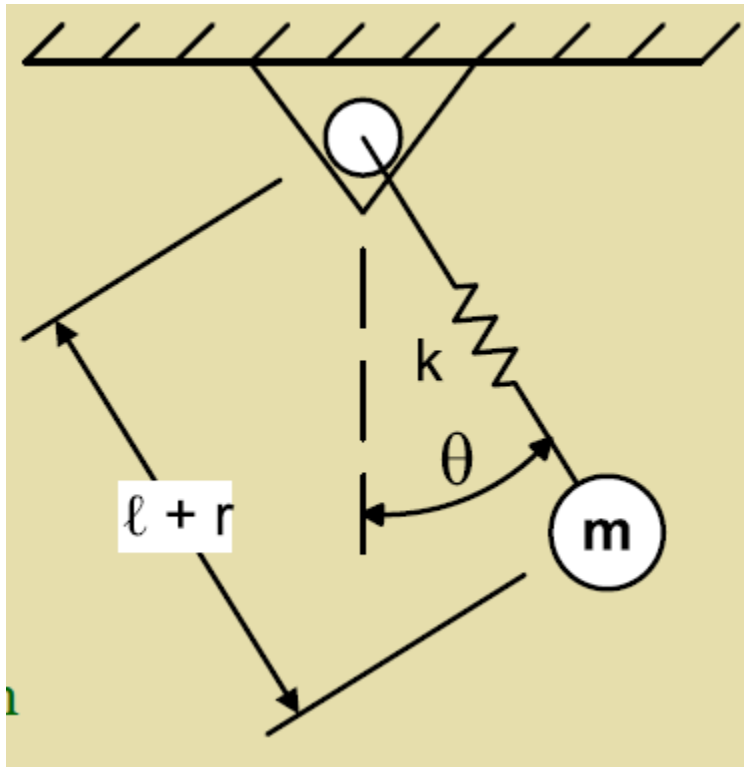
$$-ml\left(\frac{d\theta}{dt}\right)^2 = mg\cos\theta - T$$

$$ml\frac{d^2\theta}{dt^2} = -mg\sin\theta$$

Define  $\omega_0 = \sqrt{g/l}$  to find the solution:

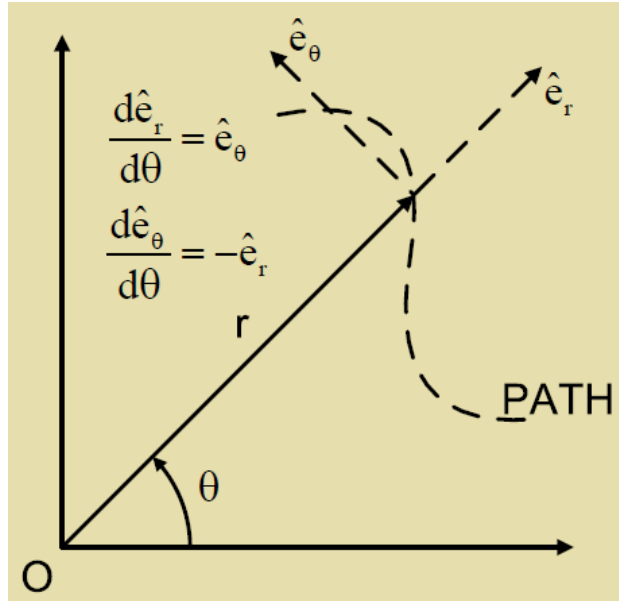
$$\frac{d^2\theta}{dt^2} = -\frac{g}{l}\sin\theta = -\omega_0^2\sin\theta$$

# Derivation of Equations of Motion



- $m$  = pendulum mass
- $m_{\text{spring}}$  = spring mass
- $\ell$  = unstretched spring length
- $k$  = spring constant
- $g$  = acceleration due to gravity
- $F_t$  = pre-tension of spring
- $r_s$  = static spring stretch,  $r_s = \frac{mg - F_t}{k}$
- $r_d$  = dynamic spring stretch
- $r$  = total spring stretch  $r_s + r_d$

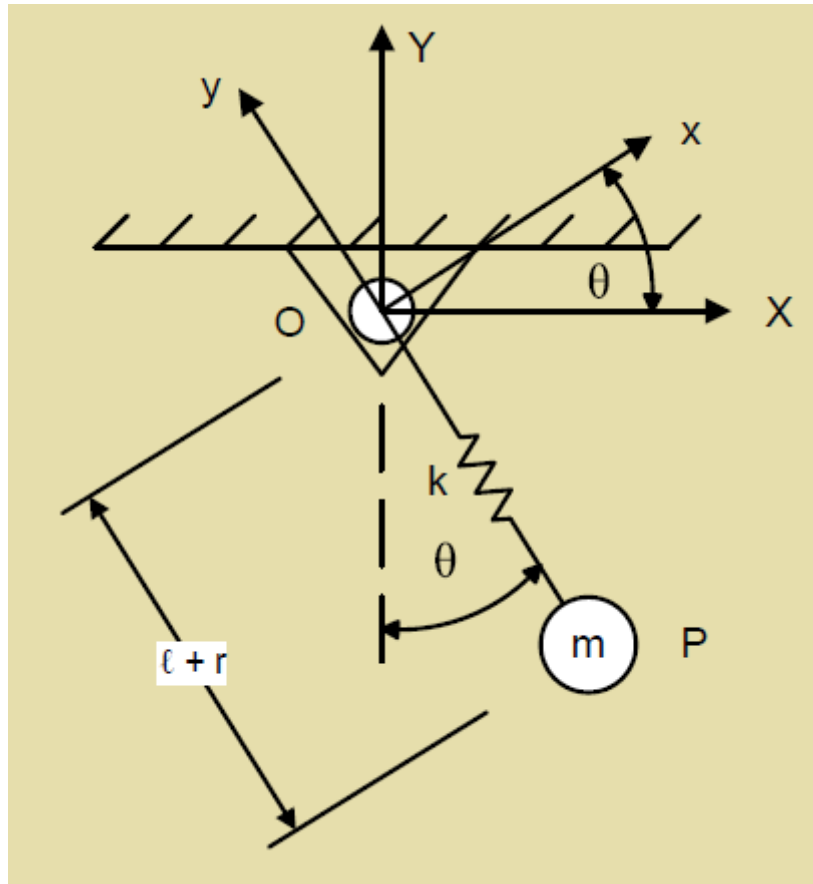
# Derivation of Equations of Motion -Polar Coordinates



- $\vec{r} = r\hat{e}_r$
- $\vec{v} = \frac{d\vec{r}}{dt} = \dot{r}\hat{e}_r + r\dot{\theta}\hat{e}_\theta = v_r\hat{e}_r + v_\theta\hat{e}_\theta$
- $\vec{a} = \frac{d\vec{v}}{dt} = (\ddot{r} - r\dot{\theta}^2)\hat{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta})\hat{e}_\theta + a_r\hat{e}_r + a_\theta\hat{e}_\theta$
- $v_r \begin{cases} \text{magnitude change} & \ddot{r} \\ \text{direction change} & \dot{r}\dot{\theta} \end{cases}$
- $v_\theta \begin{cases} \text{magnitude change} & r\ddot{\theta} + \dot{r}\dot{\theta} \\ \text{direction change} & r\dot{\theta}^2 \end{cases}$

# Derivation of Equations of Motion

## -Rigid Body Kinematics



$$\begin{bmatrix} X \\ y \\ z \end{bmatrix} = \begin{bmatrix} \cos\theta & \sin\theta & 0 \\ -\sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} X \\ Y \\ Z \end{bmatrix}$$

$$\begin{bmatrix} \hat{i} \\ \hat{j} \\ \hat{k} \end{bmatrix} = \begin{bmatrix} \cos\theta & \sin\theta & 0 \\ -\sin\theta & \cos\theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{I} \\ \hat{J} \\ \hat{K} \end{bmatrix}$$



# Derivation of Equations of Motion

## -Rigid Body Kinematics

$${}^R\vec{a}^P = {}^R\vec{a}^O + \left[ {}^R\vec{\omega}^{R_1} \times \left( {}^R\vec{\omega}^{R_1} \times \vec{r}^{OP} \right) \right] + \left[ {}^R\vec{\alpha}^{R_1} \times \vec{r}^{OP} \right] + {}^{R_1}\vec{a}^P + 2 \left[ {}^R\vec{\omega}^{R_1} \times {}^{R_1}\vec{v}^P \right]$$

$${}^R\vec{a}^O = 0$$

$${}^R\vec{\omega}^{R_1} = \dot{\theta}\hat{k} = \dot{\theta}\hat{K}$$

$$\vec{r}^{OP} = -(\ell + r)\hat{j} = -(\ell + r)\left[-\sin\theta\hat{I} + \cos\theta\hat{J}\right]$$

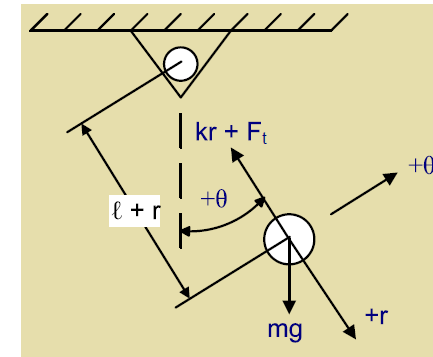
$${}^R\vec{\alpha}^{R_1} = \ddot{\theta}\hat{k} = \ddot{\theta}\hat{K}$$

$${}^{R_1}\vec{v}^P = -\dot{r}\hat{j} = -\dot{r}\left[-\sin\theta\hat{I} + \cos\theta\hat{J}\right]$$

$${}^{R_1}\vec{a}^P = -\ddot{r}\hat{j} = -\ddot{r}\left[-\sin\theta\hat{I} + \cos\theta\hat{J}\right]$$

After substitutions and evaluation:

$${}^R\vec{a}^P = \hat{i}\left[(\ell + r)\ddot{\theta} + 2\dot{r}\dot{\theta}\right] + \hat{j}\left[-\ddot{r} + (\ell + r)\dot{\theta}^2\right]$$



Free Body Diagram

$$\sum F_r = ma_r = m\left[\ddot{r} - (\ell + r)\dot{\theta}^2\right]$$

$$\sum F_\theta = ma_\theta = m\left[(\ell + r)\ddot{\theta} + 2\dot{r}\dot{\theta}\right]$$

$$m\ddot{r} - m(\ell + r)\dot{\theta}^2 + kr + F_t - mg\cos\theta = 0$$

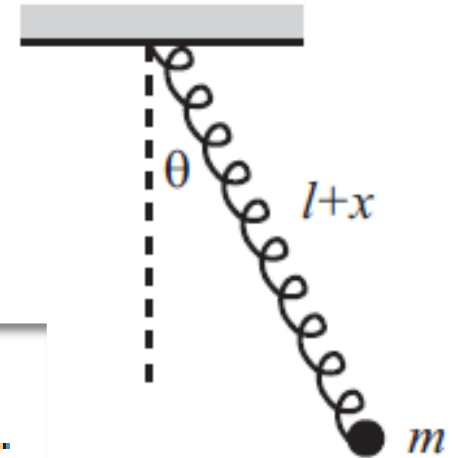
$$(\ell + r)\ddot{\theta} + 2\dot{r}\dot{\theta} + g\sin\theta = 0$$

# Derivation of Equations of Motion -Lagrange Equations

$$T = \frac{1}{2}m\left(\dot{x}^2 + (\ell + x)^2\dot{\theta}^2\right). \quad \text{Kinetic Energy}$$

$$V(x, \theta) = -mg(\ell + x)\cos\theta + \frac{1}{2}kx^2. \quad \text{Potential Energy}$$

$$L \equiv T - V = \frac{1}{2}m\left(\dot{x}^2 + (\ell + x)^2\dot{\theta}^2\right) + mg(\ell + x)\cos\theta - \frac{1}{2}kx^2.$$



# Derivation of Equations of Motion -Lagrange Equations

- Lagrange's Equation, Nonlinear equations of motion

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) = \frac{\partial L}{\partial x} \quad \Rightarrow \quad m\ddot{x} = m(\ell + x)\dot{\theta}^2 + mg \cos \theta - kx,$$

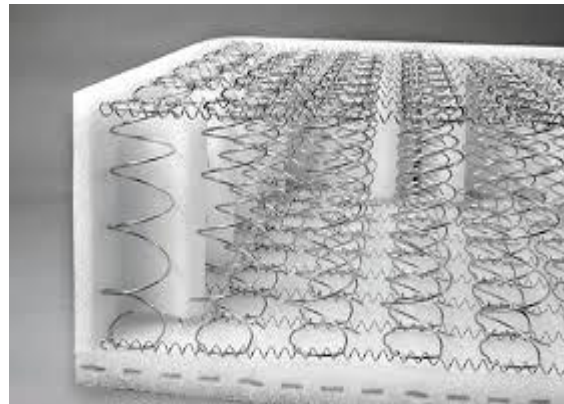
$$\begin{aligned} \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) &= \frac{\partial L}{\partial \theta} &\Rightarrow &\frac{d}{dt} (m(\ell + x)^2 \dot{\theta}) = -mg(\ell + x) \sin \theta \\ &&\Rightarrow &m(\ell + x)^2 \ddot{\theta} + 2m(\ell + x)\dot{x}\dot{\theta} = -mg(\ell + x) \sin \theta. \\ &&\Rightarrow &m(\ell + x)\ddot{\theta} + 2m\dot{x}\dot{\theta} = -mg \sin \theta. \end{aligned}$$

# Elastic pendulum in the real world

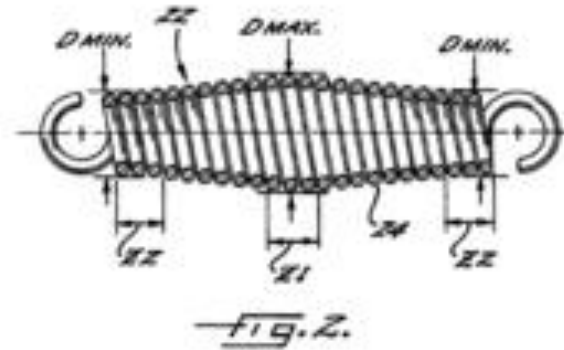
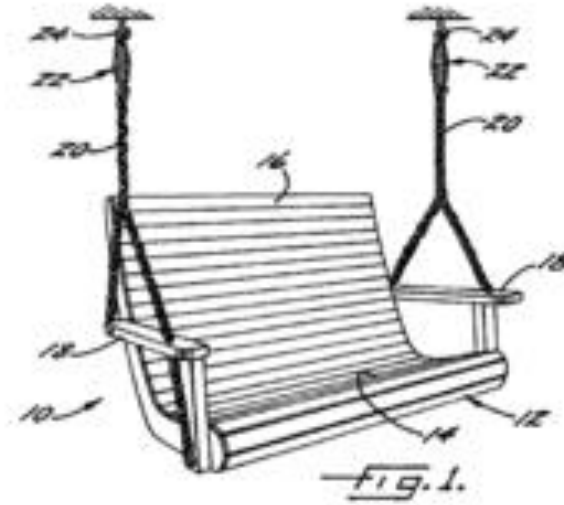
Pendulum... but not elastic:



Elastic... but not pendulum:



# Elastic pendulum in the real world – Spring Swinging



# Elastic pendulum in the real world

## -Bungee Jumping



# Trivial Cases

- System not integrable
- Initial condition without elastic potential
- Only vertical oscillation
- Initial condition with elastic potential

# Equilibrium States

- Hook's Law  $F_e = -k(l - l_0)$
- Gravitational Force  $F_g = mg$

- At equilibrium  $\sum F = F_e + F_g = 0$   
 $k(l - l_0) = mg$

- System at equilibrium

$$\ddot{x} = 0 \Rightarrow x = 0$$

$$\ddot{y} = 0 \Rightarrow y = 0$$

$$\ddot{z} = 0 \Rightarrow -\frac{k}{m} \left( \frac{l - l_0}{l} \right) z - g = 0 \Rightarrow z = -l$$

Stable state  $(x, y, z) = (0, 0, -l)$

Unstable at  $(x, y, z) = (0, 0, 2l_0 - l)$



Turning things into a handy system:

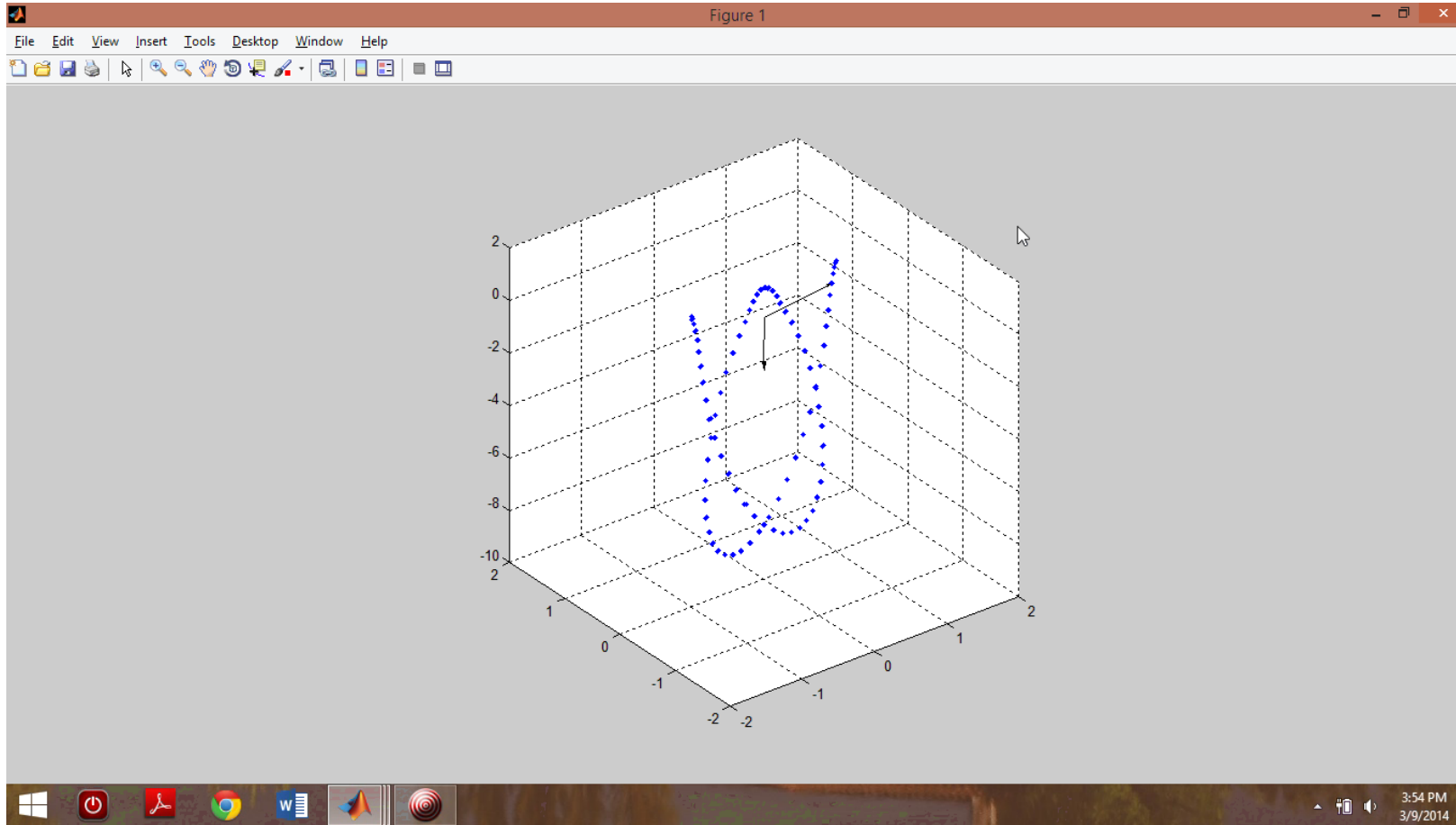
$$\ddot{x} = -\frac{\omega_z^2(r - l_0)}{r}x$$

$$\ddot{y} = -\frac{\omega_z^2(r - l_0)}{r}y$$

$$\ddot{z} = -\frac{\omega_z^2(r - l_0)}{r}z - g$$

Where  $\omega_z = \sqrt{\frac{k}{m}}$  and  $r = \sqrt{x^2 + y^2 + z^2}$

# Some regimes make familiar shapes!



Initial Conditions:

$$X = 1; V_x = 0;$$

$$Y = 0; V_y = 0;$$

$$Z = 1.1; V_z = 0;$$

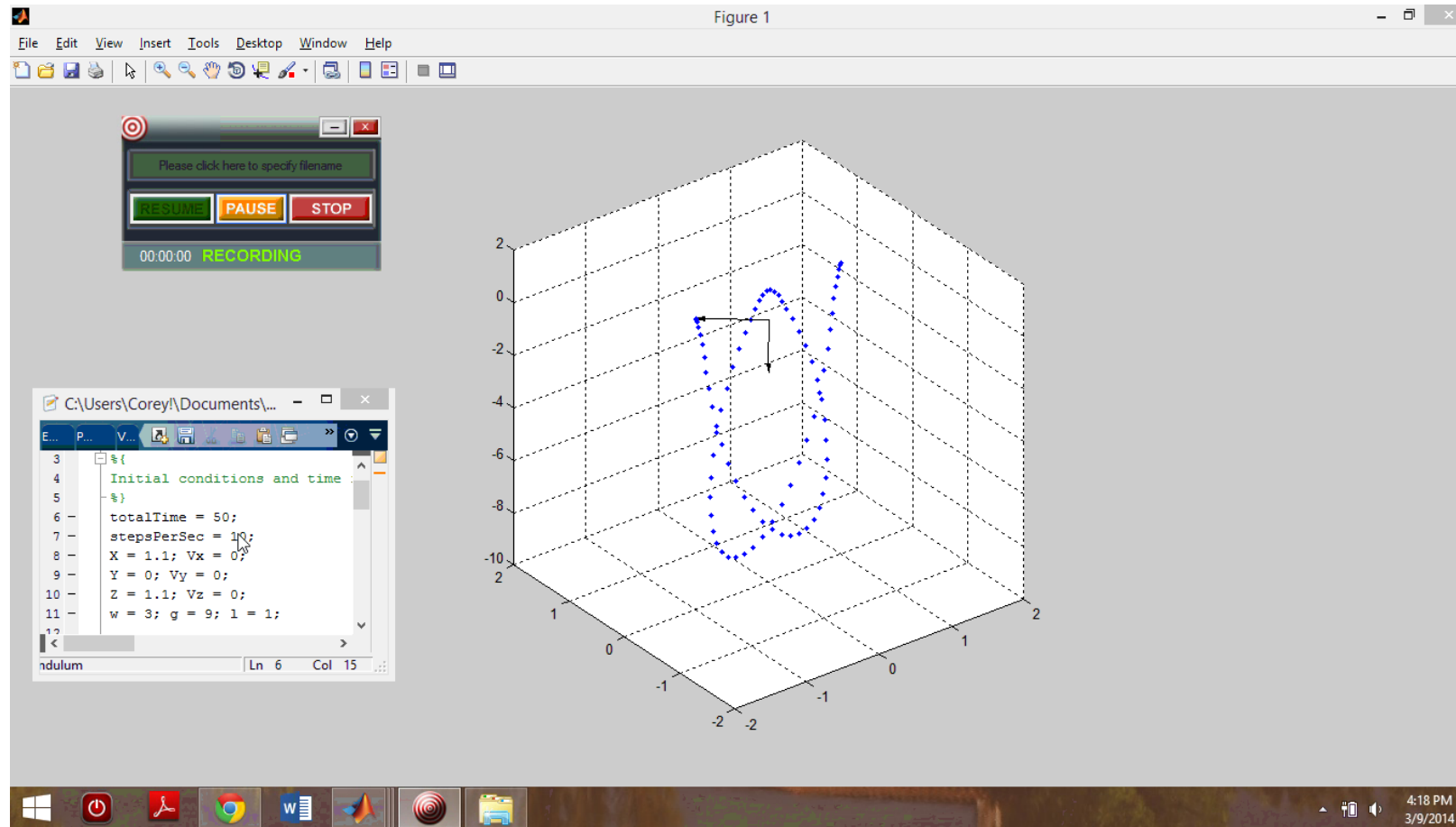
Parameters:

$$w = 3; g = 9; l = 1;$$

So motion stays in the XZ plane  
Positive Z is vertical.

Motion is shown to be relatively  
changeless over 50s

# What happens if we shake things up?



Initial Conditions

$X = 1.1$ ;  $V_x = 0$ ;

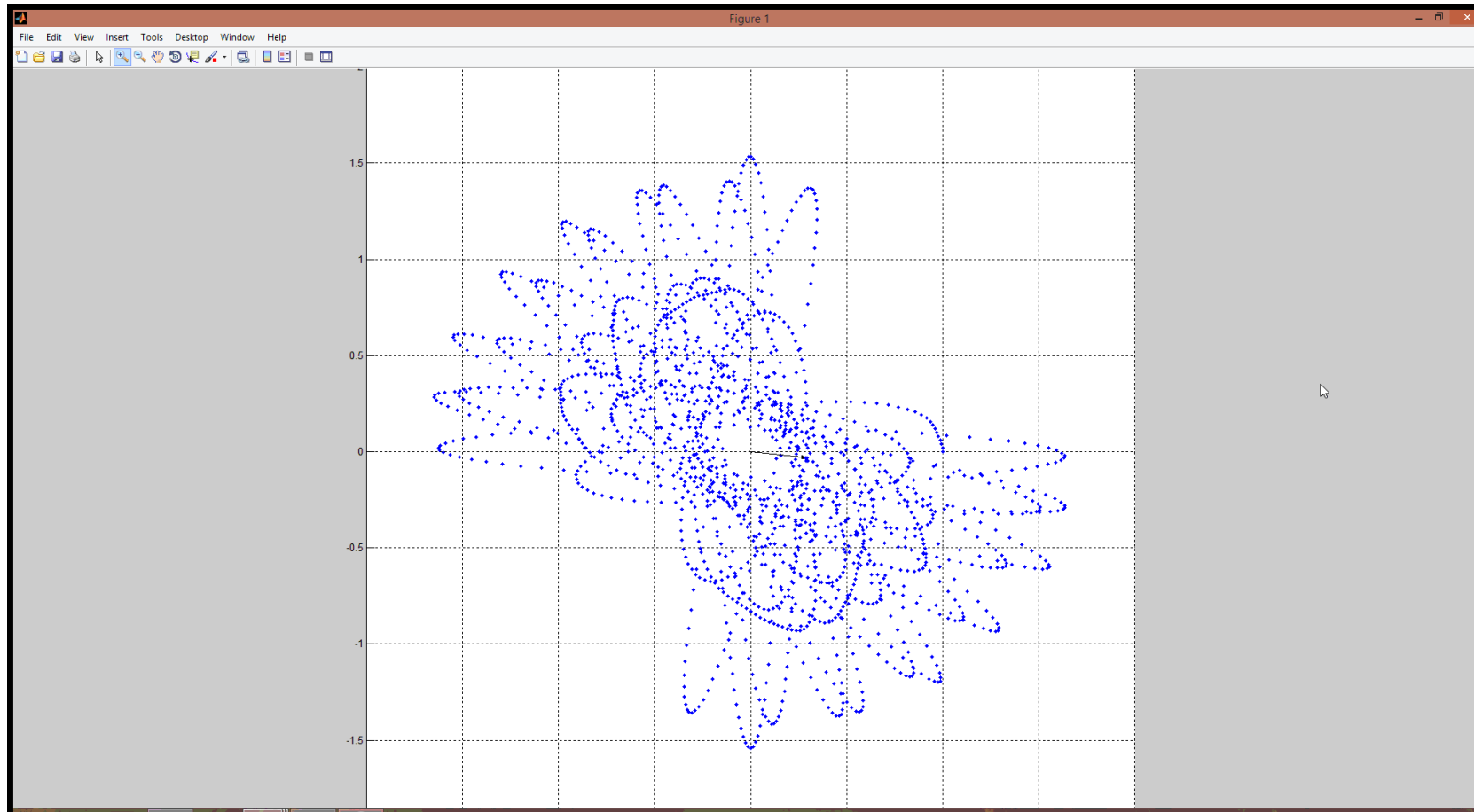
$Y = 0$ ;  $V_y = 0$ ;

$Z = 1.1$ ;  $V_z = 0$ ;

Parameters:

$w = 3$ ;  $g = 9$ ;  $l = 1$ ;

# Turning things up a bit...



Initial Conditions:

$X = 1; V_x = 0;$

$Y = 0; V_y = .2;$

$Z = 1.1; V_z = 0;$

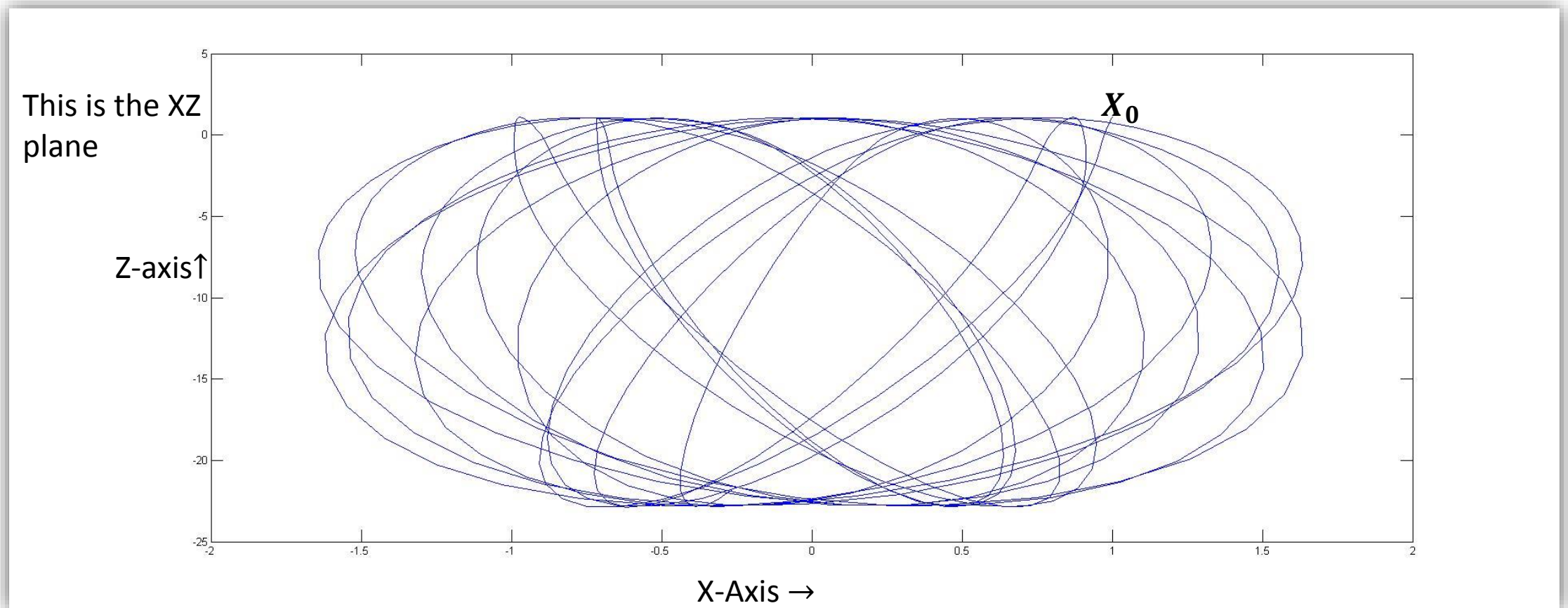
Parameters

$w = 1; g = 10; l = 1;$

$\text{totalTime} = 200;$

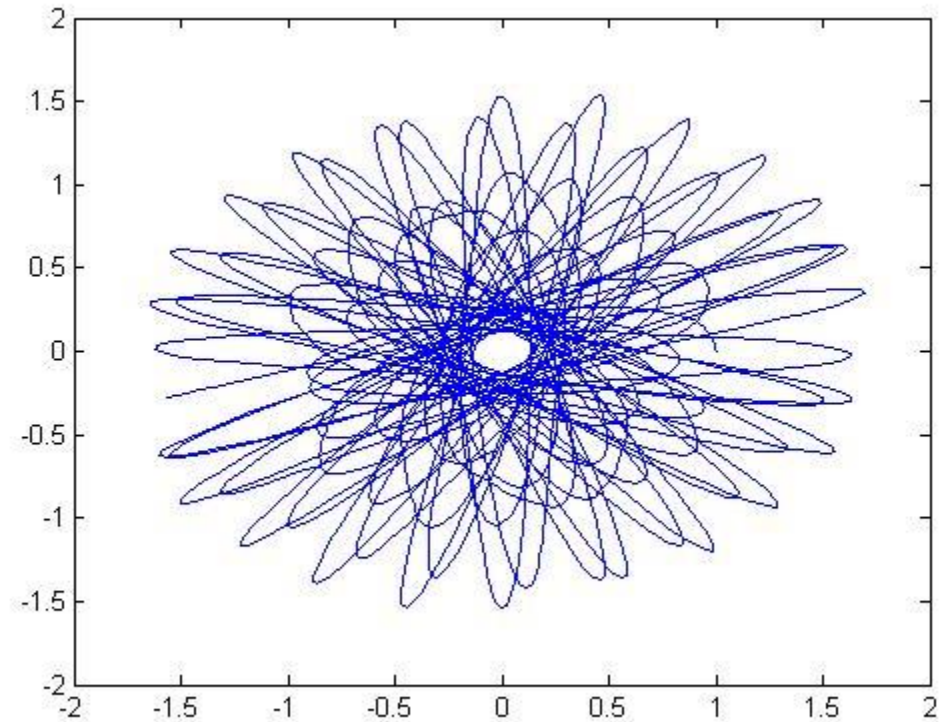
$\text{stepsPerSec} = 10;$

Let's take a closer look at the same regime:

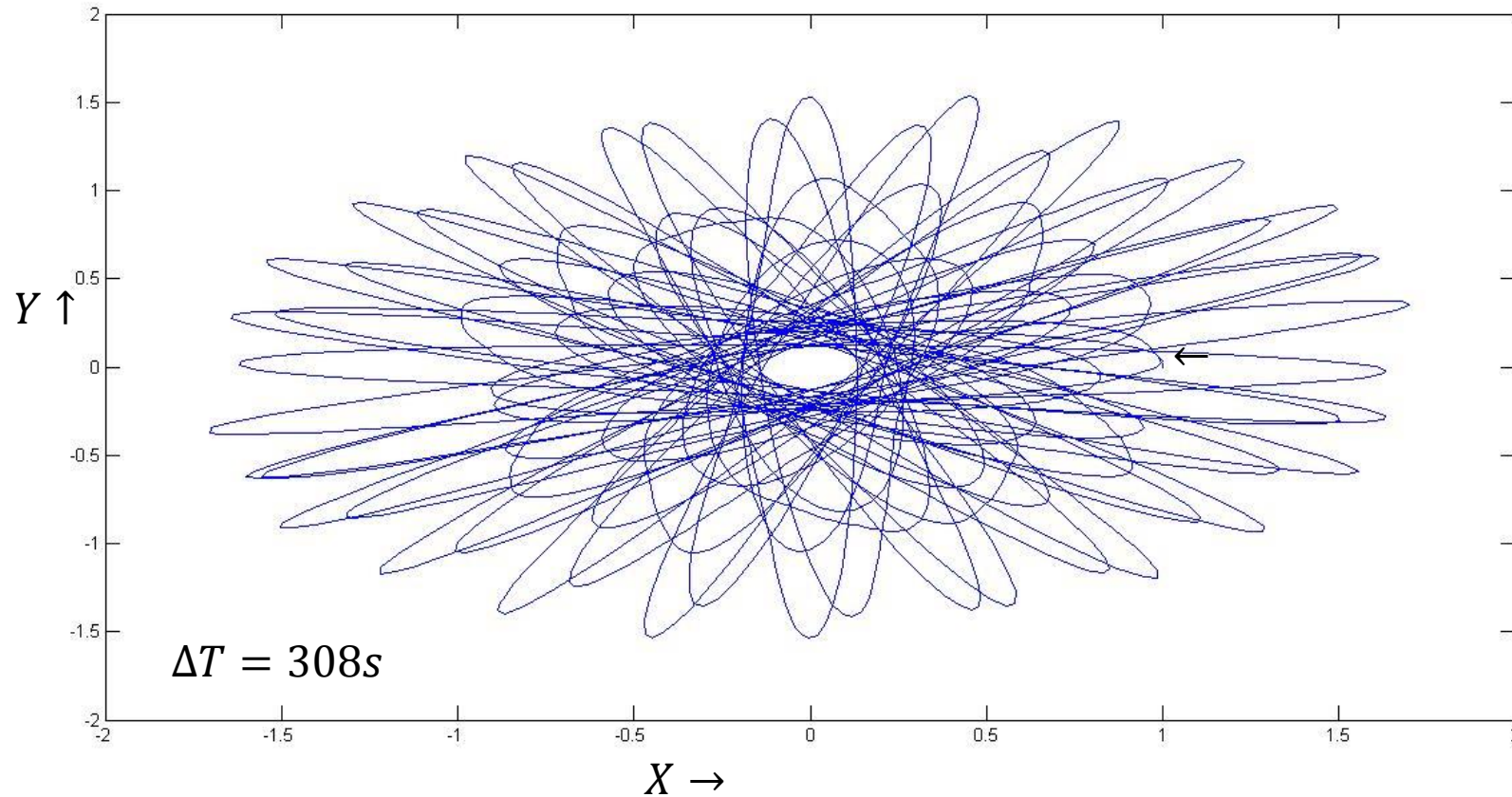


# Now look at the XY plane again.

Is it the swivel that is causing the pendulum to avoid the center?



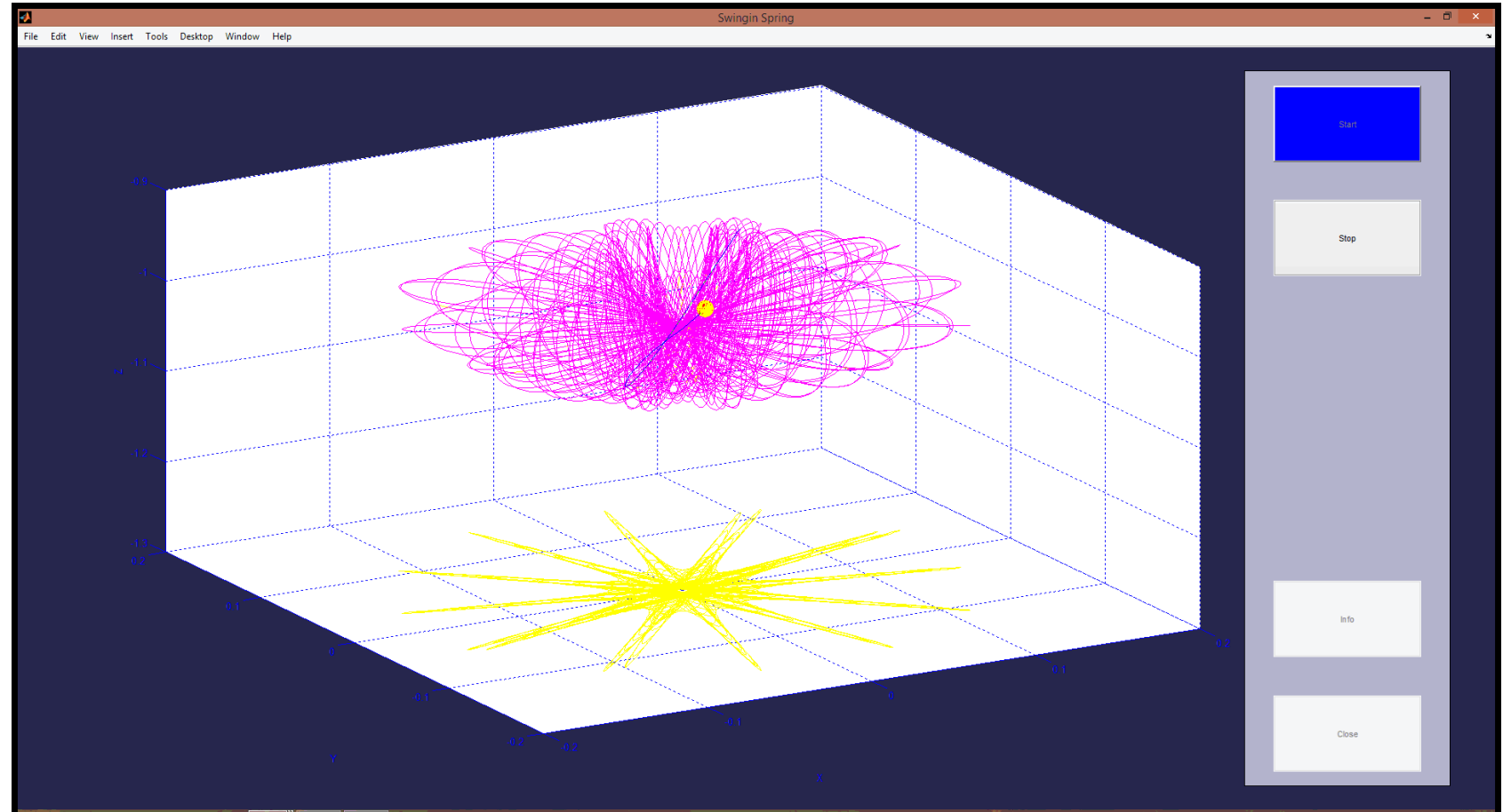
Awesome they almost meet!  
(Quasi-awesome)



# Dr. Peter Lynch's model:

Initial Conditions:

$x_0=0.01;$   
 $\dot{x}_0=0.00;$   
 $y_0=0.00;$   
 $\dot{y}_0=0.02;$   
 $z_{\text{prime}}_0=0.1;$   
 $\dot{z}_0=0.00;$





# References (1/2)

- Thanks to our mentor Joseph Gibney for getting us started on the MATLAB program and the derivations of equations of motion.
- Special thanks to Dr. Peter Lynch of the University College Dublin, Director of the UCD Meteorology & Climate Centre, for emailing his M-file and allowing us to include video of it's display of the fast oscillations of the dynamic pendulum!
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