Recall for domains  $\Omega = \mathbb{R}^2, \mathbb{R}^3$ , have "free space" Green's functions for Poisson equation

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Resolution: Use free space Green's functions as particular solutions, or use them in conjunction with symmetric reflections

# Arbitrary bounded domains

Consider Poisson equation for  $u: \Omega \to \mathbb{R}$ 

$$\Delta u = f(x, y), \quad u = 0 \text{ on } \partial \Omega$$

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Idea: use  $G_2$  as a particular solution, and write  $G(\mathbf{x}, \mathbf{x}_0) = G_2(\mathbf{x}; \mathbf{x}_0) + G_R(\mathbf{x}, \mathbf{x}_0)$ . Then  $G_R$  solves

$$\Delta_x G_R = 0, \quad G_R(\mathbf{x}, \mathbf{x}_0) = -G_2(\mathbf{x}; \mathbf{x}_0) \text{ when } \mathbf{x} \in \partial \Omega.$$

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The function  $G_R$  is the *regular part* of the Green's function, and has no singular behavior.

Main idea: If domain has certain symmetry, can use free space Green's functions as building blocks.

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Fundamental observation: for continuously differentiable  $f(x) : \mathbb{R} \to \mathbb{R}$ ,

g(x) = f(x) - f(-x) is an odd function and g(0) = 0, h(x) = f(x) + f(-x) is an even function and h'(0) = 0. Main idea: If domain has certain symmetry, can use free space Green's functions as building blocks.

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Therefore, sums/differences of free space Green's functions and their reflections can get boundary conditions right!

### Example 1: upper half space

Consider problem for  $u : \mathbb{R}^3 \cap \{z > 0\} \to \mathbb{R}$ 

$$\Delta u = f$$
,  $\lim_{|\mathbf{x}| \to \infty} u(\mathbf{x}) = 0$ ,  $u(x, y, 0) = 0$ .

Can't use free space Green's function

$$G_3(x, y, z; x_0, y_0, z_0) = -\frac{1}{4\pi\sqrt{(x-x_0)^2 + (y-y_0)^2 + (z-z_0)^2}}$$

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Using symmetry, however, suggests subtracting *G* from its mirror image about z = 0:

$$\begin{aligned} G(x, y, z; x_0, y_0, z_0) &= G_3(x, y, z; x_0, y_0, z_0) - G_3(x, y, -z; x_0, y_0, z_0) = \\ \frac{1}{4\pi} \left( \frac{1}{\sqrt{(x - x_0)^2 + (y - y_0)^2 + (z - z_0)^2}} - \frac{1}{\sqrt{(x - x_0)^2 + (y - y_0)^2 + (z + z_0)^2}} \right) \end{aligned}$$

Check that

$$\begin{aligned} G(x,y,z;x_0,y_0,z_0) &= \\ \frac{1}{4\pi} \left( \frac{1}{\sqrt{(x-x_0)^2 + (y-y_0)^2 + (z-z_0)^2}} - \frac{1}{\sqrt{(x-x_0)^2 + (y-y_0)^2 + (z+z_0)^2}} \right). \end{aligned}$$

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Equation formally satisfied by G is

$$\Delta_{x}G(\mathbf{x},\mathbf{x}_{0}) = \delta(\mathbf{x}-\mathbf{x}_{0}) - \delta(\mathbf{x}-\mathbf{x}_{0}^{*}),$$

where  $\mathbf{x}_0^* = (x_0, y_0, -z_0)$  is called the *image source*.

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Extra  $\delta$ -function is not a problem, since equation for *G* only needs to be satisfied in problem domain

$$\Delta_x G(\mathbf{x}, \mathbf{x}_0) = \delta(\mathbf{x} - \mathbf{x}_0), \quad \text{for all } \mathbf{x} \in \mathbb{R}^3 \cap \{z > 0\} \;.$$

Want (in polar coordinates)  $G(r, \theta; r_0, \theta_0)$  to solve  $\Delta G = \delta(\mathbf{x} - \mathbf{x}_0)$  with boundary condition  $G(a, \theta; r_0, \theta_0) = 0$ .

Using free space G-function

$$G_2(r, \theta; r_o, \theta_0) = \frac{1}{4\pi} \ln(r^2 + r_0^2 - 2rr_0 \cos(\theta - \theta_0)).$$

Idea is to subtract  $G_2$  from its "reflection" across the boundary, using image source at  $\mathbf{x}_0^* = a^2 \mathbf{x}_0 / r_0^2$ . The difference on the boundary is not zero, but is a constant  $= -\ln(a/r_0)/(2\pi)$ .

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$$egin{aligned} G(r, heta;r_0, heta_0) &= G_2(\mathbf{x},\mathbf{x}_0) - G_2(\mathbf{x},\mathbf{x}_0^*) + rac{1}{2\pi}\ln(a/r_0) \ &= rac{1}{4\pi}\ln\left(rac{a^2}{r_0^2}rac{r^2+r_0^2-2rr_0\cos( heta- heta_0)}{r^2+a^4/r_0^2-2ra^2/r_0\cos( heta- heta_0)}
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# Example 2: disk,cont.

Check that

$$\begin{aligned} G(r,\theta;r_0,\theta_0) &= \frac{1}{4\pi} \ln\left(\frac{a^2}{r_0^2} \frac{r^2 + r_0^2 - 2rr_0\cos(\theta - \theta_0)}{r^2 + a^4/r_0^2 - 2ra^2/r_0\cos(\theta - \theta_0)}\right) \\ &= G_2(\mathbf{x},\mathbf{x}_0) - G_2(\mathbf{x},\mathbf{x}_0^*) + \frac{1}{2\pi}\ln(a/r_0). \end{aligned}$$

satisfies the correct conditions:

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Note  $\Delta G = \delta(\mathbf{x} - \mathbf{x}_0) - \delta(\mathbf{x} - \mathbf{x}_0^*)$ , which is which  $\Delta G = \delta(\mathbf{x} - \mathbf{x}_0)$  when restricted to the disk.

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Also, evaluating G on the boundary,

$$\begin{aligned} G(a,\theta;r_o,\theta_0) &= \frac{1}{4\pi} \ln\left(\frac{a^2}{r_0^2} \frac{a^2 + r_0^2 - 2ar_0\cos(\theta - \theta_0)}{a^2 + a^4/r_0^2 - 2a^3/r_0\cos(\theta - \theta_0)}\right) \\ &= \frac{1}{4\pi} \ln\left(\frac{a^2 + r_0^2 - 2ar_0\cos(\theta - \theta_0)}{r_0^2 + a^2 - 2ar_0\cos(\theta - \theta_0)}\right) = 0. \end{aligned}$$