

## Two-dimensional eigenvalue problems, part 2

Find  $v : D \rightarrow \mathbb{R}$

$$\Delta v + \lambda v = 0, \quad \text{plus boundary conditions.}$$

## Two-dimensional eigenvalue problems, part 2

Find  $v : D \rightarrow \mathbb{R}$

$$\Delta v + \lambda v = 0, \quad \text{plus boundary conditions.}$$

- With suitable boundary conditions, eigenvalues are real, non-negative and eigenfunctions are orthogonal

$$\langle u, v \rangle = \int_D uv \, dx.$$

## Two-dimensional eigenvalue problems, part 2

Find  $v : D \rightarrow \mathbb{R}$

$$\Delta v + \lambda v = 0, \quad \text{plus boundary conditions.}$$

- With suitable boundary conditions, eigenvalues are real, non-negative and eigenfunctions are orthogonal

$$\langle u, v \rangle = \int_D uv \, dx.$$

- Difficult to write complete solution for arbitrary domain.

## Two-dimensional eigenvalue problems, part 2

Find  $v : D \rightarrow \mathbb{R}$

$$\Delta v + \lambda v = 0, \quad \text{plus boundary conditions.}$$

- With suitable boundary conditions, eigenvalues are real, non-negative and eigenfunctions are orthogonal

$$\langle u, v \rangle = \int_D uv \, dx.$$

- Difficult to write complete solution for arbitrary domain.
- Three tractable domains: rectangle, disk, surface of a sphere.

## Eigenfunctions on the rectangle

Consider domain  $D = \{0 < x < \pi, 0 < y < \pi\}$ .

## Eigenfunctions on the rectangle

Consider domain  $D = \{0 < x < \pi, 0 < y < \pi\}$ .

Natural idea to solve  $\Delta v + \lambda v = 0$ : use separation of variables!

## Eigenfunctions on the rectangle

Consider domain  $D = \{0 < x < \pi, 0 < y < \pi\}$ .

Natural idea to solve  $\Delta v + \lambda v = 0$ : use separation of variables!

Let  $v(x, y) = X(x)Y(y)$ , plug into eigenvalue equation

$$\frac{X''}{X} + \frac{Y''}{Y} = -\lambda$$

## Eigenfunctions on the rectangle

Consider domain  $D = \{0 < x < \pi, 0 < y < \pi\}$ .

Natural idea to solve  $\Delta v + \lambda v = 0$ : use separation of variables!

Let  $v(x, y) = X(x)Y(y)$ , plug into eigenvalue equation

$$\frac{X''}{X} + \frac{Y''}{Y} = -\lambda$$

By separation principle terms on left are constants  $-\lambda_x, -\lambda_y$ , so that

$$X'' + \lambda_x X = 0, \quad Y'' + \lambda_y Y = 0, \quad \lambda = \lambda_x + \lambda_y.$$



## Eigenfunctions on the rectangle, cont.

Suppose that Dirichlet boundary conditions on  $D$  were imposed; this leads to the boundary conditions

$$X(0) = 0 = X(\pi), \quad Y(0) = 0 = Y(\pi).$$

## Eigenfunctions on the rectangle, cont.

Suppose that Dirichlet boundary conditions on  $D$  were imposed; this leads to the boundary conditions

$$X(0) = 0 = X(\pi), \quad Y(0) = 0 = Y(\pi).$$

Familiar solutions are  $X = \sin(nx)$  and  $Y = \sin(my)$  for  $n, m = 1, 2, 3, \dots$

## Eigenfunctions on the rectangle, cont.

Suppose that Dirichlet boundary conditions on  $D$  were imposed; this leads to the boundary conditions

$$X(0) = 0 = X(\pi), \quad Y(0) = 0 = Y(\pi).$$

Familiar solutions are  $X = \sin(nx)$  and  $Y = \sin(my)$  for  $n, m = 1, 2, 3, \dots$

Therefore all possible combinations are

$$v_{nm}(x, y) = \sin(nx) \sin(my), \quad \lambda_{nm} = n^2 + m^2, \quad n, m = 1, 2, 3, \dots$$

(these comprise a multidimensional Fourier series)

## Eigenfunctions on the rectangle, cont.

Suppose that Dirichlet boundary conditions on  $D$  were imposed; this leads to the boundary conditions

$$X(0) = 0 = X(\pi), \quad Y(0) = 0 = Y(\pi).$$

Familiar solutions are  $X = \sin(nx)$  and  $Y = \sin(my)$  for  $n, m = 1, 2, 3, \dots$

Therefore all possible combinations are

$$v_{nm}(x, y) = \sin(nx) \sin(my), \quad \lambda_{nm} = n^2 + m^2, \quad n, m = 1, 2, 3, \dots$$

(these comprise a multidimensional Fourier series)

Since the linear operator is self adjoint, orthogonality holds

$$\int_D v_{nm} v_{n'm'} dx dy = 0 \quad \text{unless } n = n' \text{ and } m = m'.$$

## Eigenfunctions on the rectangle, cont.

Suppose that Dirichlet boundary conditions on  $D$  were imposed; this leads to the boundary conditions

$$X(0) = 0 = X(\pi), \quad Y(0) = 0 = Y(\pi).$$

Familiar solutions are  $X = \sin(nx)$  and  $Y = \sin(my)$  for  $n, m = 1, 2, 3, \dots$

Therefore all possible combinations are

$$v_{nm}(x, y) = \sin(nx) \sin(my), \quad \lambda_{nm} = n^2 + m^2, \quad n, m = 1, 2, 3, \dots$$

(these comprise a multidimensional Fourier series)

Since the linear operator is self adjoint, orthogonality holds

$$\int_D v_{nm} v_{n'm'} dx dy = 0 \quad \text{unless } n = n' \text{ and } m = m'.$$

Many other possible boundary conditions; for example, Neumann leads to  $\nabla v \cdot \hat{\mathbf{n}} = 0$  lead to

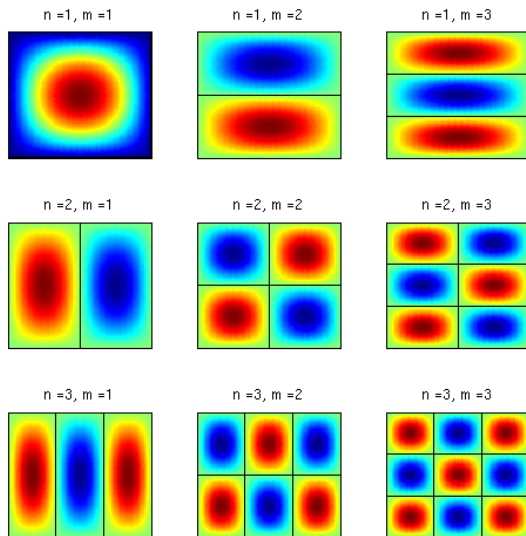
$$v_{nm}(x, y) = \cos(nx) \cos(my), \quad \lambda_{nm} = n^2 + m^2, \quad n, m = 0, 1, 2, 3, \dots$$

## Nodes of the eigenfunctions

To visualizing eigenfunctions, look at *nodes* where  $v = 0$ .

## Nodes of the eigenfunctions

To visualizing eigenfunctions, look at *nodes* where  $v = 0$ .



## Eigenfunctions on the disk

Consider circular domain  $D = \{0 < r < a, 0 < \theta < 2\pi\}$ .  
Using Laplacian in polar coordinates,

$$v_{rr} + \frac{1}{r}v_r + \frac{1}{r^2}v_{\theta\theta} = -\lambda v.$$



## Eigenfunctions on the disk

Consider circular domain  $D = \{0 < r < a, 0 < \theta < 2\pi\}$ .  
Using Laplacian in polar coordinates,

$$v_{rr} + \frac{1}{r}v_r + \frac{1}{r^2}v_{\theta\theta} = -\lambda v.$$

Look for solutions  $v = R(r)\Theta(\theta)$  and separate variables

$$\frac{r^2 R'' + rR' + \lambda r^2 R}{R} + \frac{\Theta''}{\Theta} = 0.$$

Each term must be a constant, so set  $\Theta''/\Theta = -\lambda_\theta$ .

## Eigenfunctions on the disk

Consider circular domain  $D = \{0 < r < a, 0 < \theta < 2\pi\}$ .  
Using Laplacian in polar coordinates,

$$v_{rr} + \frac{1}{r}v_r + \frac{1}{r^2}v_{\theta\theta} = -\lambda v.$$

Look for solutions  $v = R(r)\Theta(\theta)$  and separate variables

$$\frac{r^2 R'' + rR' + \lambda r^2 R}{R} + \frac{\Theta''}{\Theta} = 0.$$

Each term must be a constant, so set  $\Theta''/\Theta = -\lambda_\theta$ .

Boundary conditions for  $\Theta$  are  $2\pi$ -periodic, so  $\lambda_\theta = n^2$  and

$$\Theta = \begin{cases} \cos(n\theta), & n = 0, 1, 2, \dots \\ \sin(n\theta), & n = 1, 2, 3, \dots \end{cases}$$

Using  $\lambda_\theta = n^2$ , equation for radial part  $R(r)$  is

$$r^2 R'' + rR' + (\lambda r^2 - n^2)R = 0.$$

Using  $\lambda_\theta = n^2$ , equation for radial part  $R(r)$  is

$$r^2 R'' + rR' + (\lambda r^2 - n^2)R = 0.$$

If  $\lambda = 0$ , get Euler's equation. For the Dirichlet boundary condition, no non-zero solutions; for Neumann, get constant valued solution.

Using  $\lambda_\theta = n^2$ , equation for radial part  $R(r)$  is

$$r^2 R'' + rR' + (\lambda r^2 - n^2)R = 0.$$

If  $\lambda = 0$ , get Euler's equation. For the Dirichlet boundary condition, no non-zero solutions; for Neumann, get constant valued solution.

If  $\lambda > 0$ , simplify using change of variables  $\rho = \sqrt{\lambda}r$  so  $R = R(\rho)$  solves

$$\rho^2 R'' + \rho R' + (\rho^2 - n^2)R = 0, \quad (\text{Bessel's equation}).$$

Approximate for  $\rho$  small gives Euler's equation

$$\rho^2 R'' + \rho R' - n^2 R = 0$$

Therefore for each  $n$  expect two linearly independent solutions  $J_n(\rho), Y_n(\rho)$

$$J_n(\rho) \sim \begin{cases} \rho^n & n > 0 \\ 1 & n = 0, \end{cases} \quad Y_n(\rho) \sim \begin{cases} \rho^{-n} & n > 0 \\ \ln \rho & n = 0, \end{cases}$$

Approximate for  $\rho$  small gives Euler's equation

$$\rho^2 R'' + \rho R' - n^2 R = 0$$

Therefore for each  $n$  expect two linearly independent solutions  $J_n(\rho), Y_n(\rho)$

$$J_n(\rho) \sim \begin{cases} \rho^n & n > 0 \\ 1 & n = 0, \end{cases} \quad Y_n(\rho) \sim \begin{cases} \rho^{-n} & n > 0 \\ \ln \rho & n = 0, \end{cases}$$

- $J_n(\rho)$  and  $Y_n(\rho)$  are the *Bessel functions*.

Approximate for  $\rho$  small gives Euler's equation

$$\rho^2 R'' + \rho R' - n^2 R = 0$$

Therefore for each  $n$  expect two linearly independent solutions  $J_n(\rho), Y_n(\rho)$

$$J_n(\rho) \sim \begin{cases} \rho^n & n > 0 \\ 1 & n = 0, \end{cases} \quad Y_n(\rho) \sim \begin{cases} \rho^{-n} & n > 0 \\ \ln \rho & n = 0, \end{cases}$$

- $J_n(\rho)$  and  $Y_n(\rho)$  are the *Bessel functions*.
- Bessel functions have an infinite number of zeros; label them

$$\beta_{nm} = m\text{th positive zero of } J_n(\rho).$$



Now assume Dirichlet boundary conditions. Eigenvalues are determined by imposing boundary condition

$$R(r = a) = J_n(\sqrt{\lambda}a) = 0,$$

Now assume Dirichlet boundary conditions. Eigenvalues are determined by imposing boundary condition

$$R(r = a) = J_n(\sqrt{\lambda}a) = 0,$$

This means that for each  $n$ ,  $\sqrt{\lambda}a$  must be a zero of  $J_n$ , or

$$\lambda_{nm} = \left( \frac{\beta_{nm}}{a} \right)^2, \quad 0 \leq n < \infty, \quad 1 \leq m < \infty.$$

## Eigenfunctions on the disk, cont.

Now assume Dirichlet boundary conditions. Eigenvalues are determined by imposing boundary condition

$$R(r = a) = J_n(\sqrt{\lambda}a) = 0,$$

This means that for each  $n$ ,  $\sqrt{\lambda}a$  must be a zero of  $J_n$ , or

$$\lambda_{nm} = \left(\frac{\beta_{nm}}{a}\right)^2, \quad 0 \leq n < \infty, \quad 1 \leq m < \infty.$$

The corresponding eigenfunctions are

$$v_{nm}(x, y) = \begin{cases} J_0(\beta_{0m}r/a) & n = 0 \\ J_n(\beta_{nm}r/a) \cos(n\theta), J_n(\beta_{nm}r/a) \sin(m\theta), & n > 0. \end{cases}$$

## What *are* the Bessel functions?

To find a solution to Bessel's equation, use “method of Frobenius”:

$$R = \rho^\alpha \sum_{k=0}^{\infty} a_k \rho^k, \quad a_0 \neq 0$$

## What *are* the Bessel functions?

To find a solution to Bessel's equation, use “method of Frobenius”:

$$R = \rho^\alpha \sum_{k=0}^{\infty} a_k \rho^k, \quad a_0 \neq 0$$

Plugging in

$$\rho^\alpha \left\{ \sum_{k=0}^{\infty} [(\alpha + k)(\alpha + k - 1) + (\alpha + k) - n^2] a_k \rho^{k-2} + \sum_{k=2}^{\infty} a_{k-2} \rho^{k-2} \right\} = 0.$$

## What *are* the Bessel functions?

To find a solution to Bessel's equation, use “method of Frobenius”:

$$R = \rho^\alpha \sum_{k=0}^{\infty} a_k \rho^k, \quad a_0 \neq 0$$

Plugging in

$$\rho^\alpha \left\{ \sum_{k=0}^{\infty} [(\alpha + k)(\alpha + k - 1) + (\alpha + k) - n^2] a_k \rho^{k-2} + \sum_{k=2}^{\infty} a_{k-2} \rho^{k-2} \right\} = 0.$$

Idea: each coefficient of  $\rho^k$  must be zero.

## What *are* the Bessel functions?

To find a solution to Bessel's equation, use “method of Frobenius”:

$$R = \rho^\alpha \sum_{k=0}^{\infty} a_k \rho^k, \quad a_0 \neq 0$$

Plugging in

$$\rho^\alpha \left\{ \sum_{k=0}^{\infty} [(\alpha + k)(\alpha + k - 1) + (\alpha + k) - n^2] a_k \rho^{k-2} + \sum_{k=2}^{\infty} a_{k-2} \rho^{k-2} \right\} = 0.$$

Idea: each coefficient of  $\rho^k$  must be zero.

For  $k = 0, 1$ ,

$$[\alpha^2 - n^2]a_0 = 0, \quad [(\alpha + 1)^2 - n^2]a_1 = 0.$$

Thus  $\alpha = \pm n$  and  $a_1 = 0$  (this leads to  $a_3, a_5, a_7, \dots = 0$ ).

## What *are* the Bessel functions?

Other coefficients solve recursion relation  $a_k = -\frac{a_{k-2}}{(\alpha+k)^2 - n^2}$ ; with  $\alpha = n$  and  $a_0 = 2^{-n}/n!$ , giving

$$a_k = (-1)^k \frac{(1/2)^{2k}}{k!(n+k)!}, \quad k = 0, 2, 4, \dots$$



## What *are* the Bessel functions?

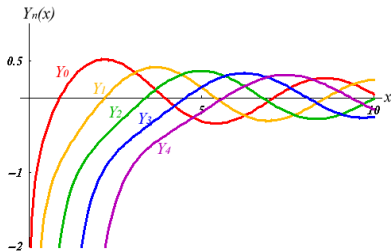
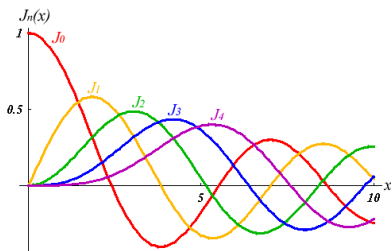
Other coefficients solve recursion relation  $a_k = -\frac{a_{k-2}}{(\alpha+k)^2 - n^2}$ ; with  $\alpha = n$  and  $a_0 = 2^{-n}/n!$ , giving

$$a_k = (-1)^k \frac{(1/2)^{2k}}{k!(n+k)!}, \quad k = 0, 2, 4, \dots$$

Finally, get convergent (for all  $\rho$ ) power series

$$J_n(\rho) = \sum_{k=0}^{\infty} (-1)^k \frac{(\rho/2)^{n+2k}}{k!(n+k)!}$$

# Bessel functions in pictures



## Orthogonality of Bessel functions

Bessel functions  $J_n(\beta_{nm}r/a)$  are themselves eigenfunctions of a differential operator,  $d^2/dr^2 + (1/r)d/dr - n^2/r^2$  ( $n$  fixed).

## Orthogonality of Bessel functions

Bessel functions  $J_n(\beta_{nm}r/a)$  are themselves eigenfunctions of a differential operator,  $d^2/dr^2 + (1/r)d/dr - n^2/r^2$  ( $n$  fixed). This operator is self-adjoint with respect to inner product

$$\langle u, v \rangle = \int_0^a u(r)v(r)rdr.$$

## Orthogonality of Bessel functions

Bessel functions  $J_n(\beta_{nm}r/a)$  are themselves eigenfunctions of a differential operator,  $d^2/dr^2 + (1/r)d/dr - n^2/r^2$  ( $n$  fixed). This operator is self-adjoint with respect to inner product

$$\langle u, v \rangle = \int_0^a u(r)v(r)rdr.$$

Thus the eigenfunctions (for each  $n$ ) are orthogonal:

$$\langle J_n(\beta_{nm}r/a), J_n(\beta_{nk}r/a) \rangle = 0, \quad m \neq k.$$

## Nodes of the eigenfunctions on a disk

$$v_{nm}(r, \theta) = \begin{cases} J_0(\beta_{0m}r/a) & n = 0 \\ J_n(\beta_{nm}r/a) \cos(n\theta), J_n(\beta_{nm}r/a) \sin(m\theta), & n > 0. \end{cases}$$

## Nodes of the eigenfunctions on a disk

$$v_{nm}(r, \theta) = \begin{cases} J_0(\beta_{0m}r/a) & n = 0 \\ J_n(\beta_{nm}r/a) \cos(n\theta), J_n(\beta_{nm}r/a) \sin(m\theta), & n > 0. \end{cases}$$

