# Asymptotic Methods

Solvability conditions and eigenvalue problems

The Fredholm alternative is a fundamental result from linear algebra, concerning the ability to solve Lx = b where L is a matrix or linear operator.

The Fredholm alternative is a fundamental result from linear algebra, concerning the ability to solve Lx = b where L is a matrix or linear operator.

Singular linear problems are frequently encountered in deriving asymptotic expansions. A useful form of the Fredholm alternative is

#### Fredholm

Suppose that  $L: V \to V$  is a linear mapping from a Hilbert space V to itself (e.g. a square matrix). If the adjoint of L has a nullspace spanned by  $v_1, v_2, \ldots, v_k$ , then necessary and sufficient conditions to solve Lv = b are

$$< v_i, b >= 0, \quad i = 1, 2, \ldots, k.$$

The Fredholm alternative is a fundamental result from linear algebra, concerning the ability to solve Lx = b where L is a matrix or linear operator.

Singular linear problems are frequently encountered in deriving asymptotic expansions. A useful form of the Fredholm alternative is

#### Fredholm

Suppose that  $L: V \to V$  is a linear mapping from a Hilbert space V to itself (e.g. a square matrix). If the adjoint of L has a nullspace spanned by  $v_1, v_2, \ldots, v_k$ , then necessary and sufficient conditions to solve Lv = b are

$$\langle v_i, b \rangle = 0, \quad i = 1, 2, ..., k.$$

Remarks:

- Proof is easy: simply take the inner product of both sides with each v<sub>i</sub>, and use the definition of adjoint.
- Notice that the solution of Lx = b will not be unique, but rather the sum of a particular solution plus the nullspace of L.
- For boundary value problems, the boundary conditions also come into play.

If A, B are nonsingular, symmetric, square matrices, then  $(A + \epsilon B)x = y$ admits a regular expansion  $x = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots$  Inserting this and equating powers of  $\epsilon$  gives a hierarchy of linear problems

$$Ax_n = -Bx_{n-1}, \quad n = 1, 2, 3, \dots$$

Suppose instead that A has a one dimensional nullspace spanned by v. Could it be that  $A + \epsilon B$  is nonsingular?

If A, B are nonsingular, symmetric, square matrices, then  $(A + \epsilon B)x = y$ admits a regular expansion  $x = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots$  Inserting this and equating powers of  $\epsilon$  gives a hierarchy of linear problems

$$Ax_n = -Bx_{n-1}, \quad n = 1, 2, 3, \dots$$

Suppose instead that A has a one dimensional nullspace spanned by v. Could it be that  $A + \epsilon B$  is nonsingular? Let's solve  $(A + \epsilon B)x = y$ . Since its unlikely that Ax = y has any solution, this choice of dominant balance is not going to work. Instead, expand  $x = \epsilon^{-1}x_{-1} + x_0 + \ldots$ , giving

$$Ax_{-1} = 0,$$
  
$$Ax_0 = y - Bx_{-1}$$

If A, B are nonsingular, symmetric, square matrices, then  $(A + \epsilon B)x = y$ admits a regular expansion  $x = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots$  Inserting this and equating powers of  $\epsilon$  gives a hierarchy of linear problems

$$Ax_n = -Bx_{n-1}, \quad n = 1, 2, 3, \dots$$

Suppose instead that A has a one dimensional nullspace spanned by v. Could it be that  $A + \epsilon B$  is nonsingular? Let's solve  $(A + \epsilon B)x = y$ . Since its unlikely that Ax = y has any solution, this choice of dominant balance is not going to work. Instead, expand  $x = \epsilon^{-1}x_{-1} + x_0 + \ldots$ , giving

$$Ax_{-1} = 0,$$
  
$$Ax_0 = y - Bx_{-1}$$

The first equation has solution  $x_{-1} = cv$  where c is yet undetermined.

If A, B are nonsingular, symmetric, square matrices, then  $(A + \epsilon B)x = y$ admits a regular expansion  $x = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots$  Inserting this and equating powers of  $\epsilon$  gives a hierarchy of linear problems

$$Ax_n = -Bx_{n-1}, \quad n = 1, 2, 3, \dots$$

Suppose instead that A has a one dimensional nullspace spanned by v. Could it be that  $A + \epsilon B$  is nonsingular? Let's solve  $(A + \epsilon B)x = y$ . Since its unlikely that Ax = y has any solution, this choice of dominant balance is not going to work. Instead, expand  $x = \epsilon^{-1}x_{-1} + x_0 + \ldots$ , giving

$$Ax_{-1} = 0,$$
  
$$Ax_0 = y - Bx_{-1}$$

The first equation has solution  $x_{-1} = cv$  where c is yet undetermined. The second equation can only be solved if the right had side is orthogonal to v, which means

$$v \cdot (y - cBv) = 0$$
, or  $c = (y \cdot v)/(Bv \cdot v)$ .

Notice this "selects" the value of c. This process can be repeated for other terms in the expansion.

Consider  $(A + \epsilon B)x = \lambda x$ , where A, B are square matrices, Suppose that  $(x_0, \lambda_0)$  is a known eigenvector/value pair for A.

Consider  $(A + \epsilon B)x = \lambda x$ , where A, B are square matrices, Suppose that  $(x_0, \lambda_0)$  is a known eigenvector/value pair for A.

Expand both  $x = x_0 + \epsilon x_1 + \ldots$  and  $\lambda = \lambda_0 + \epsilon \lambda_1 + \ldots$ , so that the leading order problem  $Ax_0 = \lambda x_0$  is automatically solved. The next order gives

$$(A - \lambda_0 I)x_1 = -Bx_0 + \lambda_1 x_0.$$

Consider  $(A + \epsilon B)x = \lambda x$ , where A, B are square matrices, Suppose that  $(x_0, \lambda_0)$  is a known eigenvector/value pair for A.

Expand both  $x = x_0 + \epsilon x_1 + \ldots$  and  $\lambda = \lambda_0 + \epsilon \lambda_1 + \ldots$ , so that the leading order problem  $Ax_0 = \lambda x_0$  is automatically solved. The next order gives

$$(A - \lambda_0 I)x_1 = -Bx_0 + \lambda_1 x_0.$$

The matrix  $A - \lambda_0 I$  is necessarily singular, and therefore so is the adjoint  $A^* - \lambda_0 I$ . Letting  $x^*$  be the adjoint eigenvector solving  $(A^* - \lambda_0 I)x^* = 0$ , Fredholm implies

$$(\lambda_1 x_0 - B x_0) \cdot x^* = 0, \quad \text{or} \quad \lambda_1 = \frac{x^* \cdot B x_0}{x^* \cdot x_0}.$$

Consider the standard Sturm-Liouville boundary value problem for u = u(x),

$$\mathcal{L}u \equiv (p(x)u')' + q(x)u = r(x), \quad u(a) = u_a, \quad u(b) = u_b.$$

If the operator  ${\cal L}$  is not singular (no zero eigenvalues), this problem has a unique solution.

Consider the standard Sturm-Liouville boundary value problem for u = u(x),

$$\mathcal{L}u \equiv (p(x)u')' + q(x)u = r(x), \quad u(a) = u_a, \quad u(b) = u_b.$$

If the operator  ${\cal L}$  is not singular (no zero eigenvalues), this problem has a unique solution.

On the other hand, if  $\mathcal{L}$  has a nullspace spanned by the eigenfunction v(x), then a solvability condition is derived by multiplying the equation by v(x) and integrating over the domain. The result is

$$\int_{a}^{b} v(x)r(x) \, dx + \left[ p(x)(u'(x)v(x) - u(x)v'(x)) \right]_{a}^{b} = 0.$$

This is both a necessary and sufficient condition to find a solution.

# Example: a perturbed SL eigenvalue problem

Consider

$$u'' + q(x)u = 0, \quad u(0) = 0 = u(L).$$

where  $q(x) = \lambda + \epsilon a(x)$ . For  $\epsilon = 0$ , this is a standard eigenvalue problem, which has solutions

$$u_0 = \sin(n\pi x/L), \quad \lambda_0 = \left(\frac{n\pi}{L}\right)^2, \quad n = 1, 2, 3, \dots$$

#### Example: a perturbed SL eigenvalue problem

Consider

$$u'' + q(x)u = 0, \quad u(0) = 0 = u(L).$$

where  $q(x) = \lambda + \epsilon a(x)$ . For  $\epsilon = 0$ , this is a standard eigenvalue problem, which has solutions

$$u_0 = \sin(n\pi x/L), \quad \lambda_0 = \left(\frac{n\pi}{L}\right)^2, \quad n = 1, 2, 3, \dots$$

We can perturb around a particular eigenpair by expanding  $\lambda = \lambda_0 + \epsilon \lambda_1 + \ldots$ and  $u = u_0 + \epsilon u_1 + \ldots$ , giving

$$\mathcal{L}u_1 \equiv u_1'' + \lambda_0 u_1 = -\lambda_1 u_0 - a(x)u_0.$$

Since  $\lambda_0$  is an eigenvalue,  $\mathcal{L}$  is singular (and self adjoint) with nullspace spanned by  $u_0$ . Solvability gives

$$\int_{a}^{b} (-\lambda_{1}u_{0} - a(x)u_{0})u_{0} dx = 0, \quad \text{or} \quad \lambda_{1} = -\frac{\int_{a}^{b} a(x)u_{0}(x)^{2} dx}{\int_{a}^{b} u_{0}(x)^{2} dx}.$$

Consider a domain which is a thin tube, described in cylindrical coordinates as  $\{(r, Z)|0 < Z < L, r < R_0f(z/L)\}$  where  $R_0 \ll L$ . Want to solve the steady state diffusion equation

$$\Delta T(r, Z) = T_{rr} + \frac{1}{r}T_r + T_{ZZ} = 0,$$
  

$$T(r, 0) = 0, \quad T(r, 1) = T_0, \quad \nabla T \cdot n = 0 \quad \text{on } r = f(Z/L).$$

Here the outward normal is proportional to  $(1, -(R_0/L)f'(z/L))$ .

Consider a domain which is a thin tube, described in cylindrical coordinates as  $\{(r, Z)|0 < Z < L, r < R_0f(z/L)\}$  where  $R_0 \ll L$ . Want to solve the steady state diffusion equation

$$\Delta T(r, Z) = T_{rr} + \frac{1}{r}T_r + T_{ZZ} = 0,$$
  

$$T(r, 0) = 0, \quad T(r, 1) = T_0, \quad \nabla T \cdot n = 0 \quad \text{on } r = f(Z/L).$$

Here the outward normal is proportional to  $(1, -(R_0/L)f'(z/L))$ .

Nondimensionalize  $u = T/T_0$ ,  $\rho = r/R_0$ , z = Z/L, giving

$$u_{\rho\rho} + \frac{1}{\rho} u_{\rho} + \epsilon^2 u_{zz} = 0, \quad u(\rho, 0) = 0, \quad u(\rho, 1) = 1, \quad u_{\rho} = \epsilon^2 f'(z) u_z \quad \text{on } \rho = f(z).$$
  
where  $\epsilon = R_0/L$ .

#### Example: diffusion through a slender body, cont.

Expanding  $u = u_0 + \epsilon u_1 + \ldots$  leads to

$$egin{aligned} \mathcal{L} u_0 &= u_{0
ho
ho} + rac{1}{
ho} = 0, \quad 0 \leq 
ho < f(z), \quad u_{0
ho}(f(z),z) = 0, \ \mathcal{L} u_1 &= -u_{0zz} = 0 \leq 
ho < f(z), \quad u_{1
ho}(f(z),z) = f'(z) u_{0z}. \end{aligned}$$

The solution for the first is  $u_0 = A(z)$ . The linear operator  $\mathcal{L}$  is self adjoint with respect to the inner product  $\int_0^{f(z)} v_1(\rho)v_2(\rho)\rho d\rho$ . The second equation therefore has a solvability condition obtained by multiplying by the eigenfunction 1 and integrating:

$$-\int_0^{f(z)} u_{0zz}\rho d\rho = \int_0^{f(z)} (\rho u_{1\rho})_\rho d\rho = \rho u_{1\rho}\Big|_0^{f(z)} = f(z)f'(z)u_{0z}.$$

#### Example: diffusion through a slender body, cont.

Expanding  $u = u_0 + \epsilon u_1 + \ldots$  leads to

$$egin{aligned} \mathcal{L} u_0 &= u_{0
ho
ho} + rac{1}{
ho} = 0, \quad 0 \leq 
ho < f(z), \quad u_{0
ho}(f(z),z) = 0, \ \mathcal{L} u_1 &= -u_{0zz} = 0 \leq 
ho < f(z), \quad u_{1
ho}(f(z),z) = f'(z) u_{0z}. \end{aligned}$$

The solution for the first is  $u_0 = A(z)$ . The linear operator  $\mathcal{L}$  is self adjoint with respect to the inner product  $\int_0^{f(z)} v_1(\rho) v_2(\rho) \rho d\rho$ . The second equation therefore has a solvability condition obtained by multiplying by the eigenfunction 1 and integrating:

$$-\int_{0}^{f(z)}u_{0zz}\rho d\rho = \int_{0}^{f(z)}(\rho u_{1\rho})_{\rho}d\rho = \rho u_{1\rho}\Big|_{0}^{f(z)} = f(z)f'(z)u_{0z}.$$

Thus  $f(z)^2 A''(z) + f(z)f'(z)A'(z) = 0$ , which can be compactly written as

$$\left[f(z)^2A_z\right]_z=0.$$

This is just one dimensional steady state diffusion with variable diffusivity. Using A(0) = 0 and A(1) = 1, one can solve by direct integration

$$A(z) = \frac{\int_0^z f(z')^2 dz'}{\int_0^1 f(z')^2 dz'}.$$